EPJ Web of Conferences **81**, 03004 (2014) DOI: 10.1051/epjconf/20148103004 © Owned by the authors, published by EDP Sciences, 2014

A direct test of time-reversal symmetry in the neutral *K* meson system with $K_S \rightarrow \pi \ell \nu$ and $K_L \rightarrow 3\pi^0$ at KLOE-2

Aleksander Gajos^{1,a} on behalf of the KLOE-2 Collaboration

¹ Institute of Physics, Jagiellonian University, ul. Reymonta 4, 30-059 Cracow, Poland

Abstract. Quantum entanglement of *K* and *B* mesons allows for a direct experimental test of time-reversal symmetry independent of CP violation. The \mathcal{T} symmetry can be probed by exchange of initial and final states in the reversible transitions between flavor and CP-definite states of the mesons which are only connected by the \mathcal{T} conjugation. While such a test was successfully performed by the BaBar experiment with neutral *B* mesons, the KLOE-2 detector can probe \mathcal{T} -violation in the neutral kaons system by investigating the process with $K_S \to \pi^{\pm} l^{\mp} v_l$ and $K_L \to 3\pi^0$ decays. Analysis of the latter is facilitated by a novel reconstruction method for the vertex of $K_L \to 3\pi^0$ decay which only involves neutral particles. Details of this new vertex reconstruction technique are presented as well as prospects for conducting the direct \mathcal{T} symmetry test at the KLOE-2 experiment.

1 Introduction

Among possible experimental ways to study the \mathcal{T} symmetry violation, it is of special interest to test the symmetry directly, i.e. by comparing amplitudes for a process and its time inverse. For spin 0 particles such as neutral mesons the inverse process is obtained simply by the exchange of initial and final states. To date, the only evidence of \mathcal{T} violation in the neutral kaon system was found by the CPLEAR experiment through measurement of the Kabir asymmetry [1]. However, use of the $C\mathcal{PT}$ -even $K^0 \subseteq \overline{K}^0$ process raised some controversy due to possible influence of $C\mathcal{P}$ violation on the result. Quantum entanglement of neutral kaons produced at the ϕ factory allows to obtain and compare kaon transitions between flavour-definite and $C\mathcal{P}$ -definite states and their time inverses which are only connected by time reversal conjugation [2]. This allows for a direct test of the \mathcal{T} symmetry independent of $C\mathcal{P}$ and $C\mathcal{PT}$. A similar idea was recently used by the BaBar experiment to directly observe \mathcal{T} violation in the neutral *B* meson system [3, 4]. In turn, KLOE-2 is capable of performing the first direct \mathcal{T} symmetry test with neutral kaons.

2 Principle of the test

For a direct \mathcal{T} symmetry test with neutral kaons, a set of transitions must be chosen such that their \mathcal{T} -inverses can be observed as well and their *in* and *out* states may be unambiguously identified by

This is an Open Access article distributed under the terms of the Creative Commons Attribution License 4.0, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

^ae-mail: aleksander.gajos@uj.edu.pl

Article available at http://www.epj-conferences.org or http://dx.doi.org/10.1051/epjconf/20148103004

EPJ Web of Conferences

observation of kaon decay final states. These conditions are met by states with definite strangeness $\{K^0, \bar{K}^0\}$ and $C\mathcal{P}$ -eigenstates $\{K_+, K_-\}$. The former are identified by semileptonic decays $K^0 \to \pi^- \ell^+ \nu_l$ and $\bar{K}^0 \to \pi^+ \ell^- \bar{\nu}_l$ (with assumption of the $\Delta Q = \Delta S$ rule) whereas the latter must decay hadronically into two pions $(\pi^+\pi^-, \pi^0\pi^0)$ for $C\mathcal{P}=+1$ or $3\pi^0$ for $C\mathcal{P}=-1$. These two bases are connected by four possible transitions, listed in Table 1. Independence of the measured asymmetry of $C\mathcal{P}$ -violating effects is guaranteed by the fact that for any transition its time inverse is not identical with its $C\mathcal{P}$ conjugate, by contrast with e.g. the Kabir asymmetry in $K^0 \subseteq \bar{K}^0$. Probability of each transition

Table 1. Transitions between flavour-definite and CP-definite states of neutral kaons and their time-reversal conjugates. Each of the transitions is experimentally identified by a time-ordered pair of kaon decays.

	Transition	Identified by	\mathcal{T} -conjugate	Identified by
1	$K^0 \to K_+$	$(\ell^-, \pi\pi)$	$K_+ \rightarrow K^0$	$(3\pi^0, \ell^+)$
2	$K^0 \rightarrow K$	$(\ell^{-}, 3\pi^{0})$	$K_{-} \rightarrow K^{0}$	$(\pi\pi,\ell^+)$
3	$\bar{K}^0 \rightarrow K_+$	$(\ell^+,\pi\pi)$	$K_+ \rightarrow \bar{K}^0$	$(3\pi^0,\ell^-)$
4	$\bar{K}^0 \to K$	$(\ell^+, 3\pi^0)$	$K_{-} \rightarrow \bar{K}^{0}$	$(\pi\pi,\ell^-)$

can be compared with its time-reversal conjugate in search of a discrepancy which would signal \mathcal{T} -violation. Experimentally, final states of kaons in the transitions would be identified directly by recording their decays while recognition of a living kaon state is uniquely possible at KLOE-2 as kaons produced in the ϕ meson decay exhibit quantum entanglement which guarantees the living kaon to be in an orthogonal state to its first-decaying partner. This way the double decay rates can be collected as a function of proper decay time difference for the two kaons decaying through chosen channels and the following ratios can be defined as observables of the \mathcal{T} symmetry test:

$$R_1^{exp}(\Delta t) = \mathbf{I}(\ell^-, \pi \pi; \Delta t) / \mathbf{I}(3\pi^0, \ell^+; \Delta t),$$
(1)

$$R_2^{exp}(\Delta t) = I(\ell^-, 3\pi^0; \Delta t) / I(\pi\pi, \ell^+; \Delta t),$$
(2)

$$R_{3}^{exp}(\Delta t) = I(\ell^{+}, \pi \pi; \Delta t) / I(3\pi^{0}, \ell^{-}; \Delta t),$$
(3)

$$R_4^{exp}(\Delta t) = I(\ell^+, 3\pi^0; \Delta t) / I(\pi\pi, \ell^-; \Delta t).$$
(4)

Among the above ratios, R_2^{exp} and R_4^{exp} concern processes for which statistics sufficient for a significant test is expected by KLOE-2 [2]. These experimental observables are related to ratios of amplitudes by the following proportionality [5]:

$$R_2(\Delta t) = P[K^0(0) \to K_-(\Delta t)] / P[K_-(0) \to K^0(\Delta t)] = R_2^{exp}(\Delta t)/C,$$
(5)

$$R_4(\Delta t) = P[\bar{K}^0(0) \to K_-(\Delta t)] / P[K_-(0) \to \bar{K}^0(\Delta t)] = R_4^{exp}(\Delta t)/C,$$
(6)

where the constant $C = \frac{BR(K_L \rightarrow 3\pi^0) \cdot \Gamma_L}{BR(K_S \rightarrow \pi\pi) \cdot \Gamma_S}$ involves kaon parameters well determined i.a. by the KLOE experiment.

After extraction of the R_2 and R_4 probability ratios from (2) and (4), their asymptotic behaviour for $\Delta t \gg \tau_S$ can be compared with the theoretical expectation:

$$R_2(\Delta t \gg \tau_s) \simeq 1 - 4\Re\epsilon, \qquad R_4(\Delta t \gg \tau_s) \simeq 1 + 4\Re\epsilon,$$
(7)

in order to measure the \mathcal{T} -violating parameter $\Re \epsilon$ [2].

3 Reconstruction of events for the test at KLOE

The KLOE detector is located at the DA Φ NE e^+e^- collider, a ϕ -factory operating at $\sqrt{s} \approx 1020$ MeV. In the years 1999–2006 KLOE has collected 2.5 fb⁻¹ of data. KLOE is a barrel-shaped detector whose MESON 2014-13th International Workshop on Production, Properties and Interaction of Mesons

basic components are large drift chamber (DC) and electromagnetic calorimeter (EMC) immersed in magnetic field of 0.52 T. Recently the detector was upgraded to KLOE-2 [6] with addition of new calorimeters at small angles around the beam pipe [7] and a new Cylindrical-GEM inner tracker [8]. Processes required for the \mathcal{T} test include semileptonic decays of neutral kaons with the partner kaon decaying into 2 or 3 pions. While for the 2-pion final state $\pi^+\pi^-$ can be chosen and well reconstructed from DC tracks, the $K_L \to 3\pi^0$ decay requires special treatment as it only includes neutral particles and the $K_S \to \pi \ell \nu$ decay does not provide full kinematic information on the event due to a missing neutrino. Therefore a special reconstruction method for $K_L \to 3\pi^0 \to 6\gamma$ decay was prepared which uses only information on γ hits in the EMC. The decay point and time are reconstructed using a technique similar to GPS positioning. More details can be found in Ref. [9].

Acknowledgements

This work was supported in part by the EU Integrated Infrastructure Initiative Hadron Physics Project under contract number RII3-CT- 2004-506078; by the European Commission under the 7th Framework Programme through the *Research Infrastructures* action of the *Capacities* Programme, Call: FP7-INFRASTRUCTURES-2008-1, Grant Agreement No. 227431; by the Polish National Science Centre through the Grants No. 0469/B/H03/2009/37, 0309/B/H03/2011/40, 2011/03/N/ST2/02641, 2011/01/D/ST2/00748, 2011/03/N/ST2/02652, 2013/08/M/ST2/00323 and by the Foundation for Polish Science through the MPD programme and the project HOMING PLUS BIS/2011-4/3.

References

- [1] A. Angelopoulos et al. [CPLEAR Collaboration], Phys. Lett. B 444, 43 (1998)
- [2] J. Bernabeu, A. Di Domenico and P. Villanueva-Perez, Nucl. Phys. B 868, 102 (2013)
- [3] J. Bernabeu, F. Martinez-Vidal and P. Villanueva-Perez, JHEP 1208, 64 (2012)
- [4] J. P. Lees *et al.* [BaBar Collaboration], Phys. Rev. Lett. **109**, 211801 (2012)
- [5] K. R. Schubert, L. L. Gioi, A. J. Bevan and A. Di Domenico, arXiv:1401.6938 [hep-ex].
- [6] D. Moricciani [KLOE-2 Collaboration], PoS EPS -HEP2011, 198 (2011)
- [7] D. Domenici, PoS EPS -HEP2013, 495 (2014)
- [8] A. Balla, G. Bencivenni, P. Branchini et al., Nucl. Instrum. Meth. A 732, 221 (2013)
- [9] A. Gajos, arXiv:1409.2132 [hep-ex]