

## MEASUREMENT OF THE $\eta \rightarrow \pi^+\pi^-\pi^0$ DECAY WITH WASA-at-COSY DETECTOR

PATRIK ADLARSON

*Institute for Physics and Astronomy, Uppsala University, Uppsala, 751 20, Sweden*  
*patrik.adlarson@fysast.uu.se*

MARCIN ZIELIŃSKI

*Institute of Physics, Jagiellonian University Kraków, 30-059 Kraków, Poland*  
*m.zielinski@uj.edu.pl*

FOR THE WASA-at-COSY COLLABORATION

One of the objectives of the physics programme of the WASA-at-COSY facility is to study the isospin violating  $\eta$  hadronic decays into  $\pi^+\pi^-\pi^0$  systems driven by the term of QCD Lagrangian which depends on the d and u quark mass difference. These studies can be made in terms of the Dalitz plot parameters describing the density population which is proportional to the square of the amplitude  $|A(x, y)|^2$ . This contribution describes the current status of the analysis of the  $\eta \rightarrow \pi^+\pi^-\pi^0$  decay in the  $pd \rightarrow {}^3\text{He}\eta$  and as well in the  $pp \rightarrow pp\eta$  reaction with WASA-at-COSY.

*Keywords:* Meson production; hadronic decays; ChPT; WASA-at-COSY.

PACS numbers: 11.25.Hf, 123.1K

### 1. Motivation

The isospin violating strong decay  $\eta \rightarrow \pi^+\pi^-\pi^0$  allows access to light quark mass ratios. At lowest order of chiral perturbation theory (ChPT) the amplitude is proportional to the light quark mass difference ( $m_d - m_u$ ) and may be written as

$$A \propto \frac{m_d - m_u}{F_\pi^2} \left( 1 + \frac{3(s - s_0)}{m_\eta^2 - m_\pi^2} \right), \quad (1)$$

where  $F_\pi$  is the pion decay constant,  $s = (p_{\pi^+} + p_{\pi^-})^2 = (p_\eta - p_{\pi^0})^2$  and  $s_0 = \frac{1}{3}(m_\eta^2 + 2m_{\pi^+}^2 + m_{\pi^0}^2)$ . At higher order of ChPT it has been found that final state pion interaction contribute to the decay width.<sup>1,2</sup> The decay width scales as  $\Gamma = \left(\frac{Q_D}{Q}\right)^4 \bar{\Gamma}$ , where  $Q^2 = \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2}$ ,  $\hat{m} = \frac{1}{2}(m_u + m_d)$ , and the decay width  $\bar{\Gamma}$  and  $Q_D = 24.2$  are calculated in the Dashen limit.<sup>3</sup> This scaling works under the pre-requisite that  $\bar{\Gamma}$  is understood reliably. To test this, theoretical predictions and experimental measurements of pion kinematical distributions may be compared in a Dalitz plot, where the axes are defined as  $x = \sqrt{3}\frac{T_+ - T_-}{Q_\eta}$ ,  $y = \frac{3T_0}{Q_\eta} - 1$ . Here  $T_+$ ,

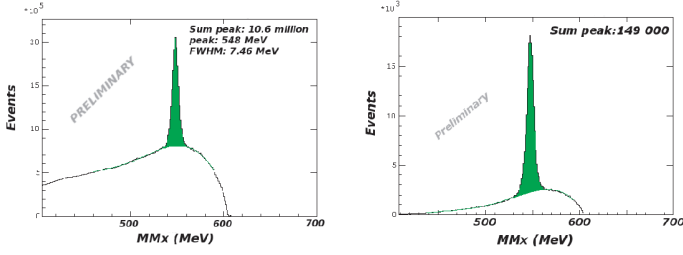


Fig. 1. (left) Missing mass for the 2008 data calculated from the identified  ${}^3\text{He}$ . (right) Missing mass after selecting  $3\pi$  candidates and including a cut on  $MM({}^3\text{He}\pi^+\pi^-)$  and  $MM({}^3\text{He}\pi^0)$ .

$T_-$  and  $T_0$  denote the kinetic energies of pions in the rest frame of the  $\eta$  meson, and  $Q_\eta = T_+ + T_- + T_0 = m_\eta - 2m_{\pi^+} - m_{\pi^0}$ . The standard way to parametrize the Dalitz plot density is a polynomial expansion around the center point:  $|A(x, y)|^2 \propto 1 + ay + by^2 + dx^2 + fy^3 + \dots$  where  $a, b, d, f$  are the Dalitz plot parameters. The experimental results are dominated by KLOE with a Dalitz plot containing  $1.34 \cdot 10^6$  events.<sup>4</sup> This result shows a significant deviation of parameters  $b$  and  $f$  in comparison to the theoretical predictions based on ChPT. It is therefore important to perform an independent measurement, which is one of the aims of the WASA-at-COSY.

## 2. $pd \rightarrow {}^3\text{He}\eta$ measurement

In 2008 and 2009 WASA-at-COSY<sup>5</sup> measured  $pd \rightarrow {}^3\text{He}X$  reaction at beam energy 1 GeV, collecting  $10^7$  and  $2 \cdot 10^7$   $\eta$  mesons respectively. The missing mass with respect to  ${}^3\text{He}$  is used to tag the  $\eta$  meson (Fig. 1 left). In addition two tracks of opposite charge are required in the Mini Drift Chamber in the angular range  $30.5^\circ < \theta < 150^\circ$ . Furthermore two  $\gamma$  with an invariant mass close to  $\pi^0$  are required. The  $pd \rightarrow {}^3\text{He}\pi\pi$  reaction is reduced by imposing conditions on the missing mass calculated for  ${}^3\text{He}\pi^+\pi^-$  and the missing mass calculated for  ${}^3\text{He}\pi^0$ . The preliminary analysis yields 149 000  $\eta \rightarrow \pi^+\pi^-\pi^0$  candidates from the 2008 data, shown in Fig. 1 right. The experimental resolution is better for the  $\eta$  four-momenta from  ${}^3\text{He}$  compared to the information derived from the  $\eta$  decay products. Therefore a kinematical fit for the reaction  $pd \rightarrow {}^3\text{He}\pi^+\pi^-\pi^0$  has been used with  ${}^3\text{He}$  observables fixed and a cut on the 1% level of the probability density function. To estimate the  $\eta$  content in each Dalitz plot bin, a four-degree polynomial fit is performed over the background region. The preliminary experimental results for the  $x, y$  projections of the Dalitz plot are compared in Fig. 2 to Monte Carlo simulations of the  $\eta \rightarrow \pi^+\pi^-\pi^0$  weighted with the tree-level prediction (Eq. (1)).

## 3. $pp \rightarrow pp\eta$ measurement

The measurement of the  $pp \rightarrow ppX$  reaction was conducted in 2008 and in 2010 at beam kinetic energy 1.4 GeV. The collected sample of data yields about  $10^8$  produced  $\eta$  mesons. Protons and charged pions were detected using scintillators

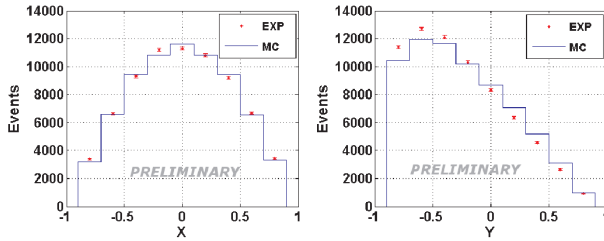


Fig. 2. Projections of Dalitz Plot, not corrected for acceptance and normalized to sum of experimental data: **(left)** X-projection **(right)** Y-projection. Solid line indicates MC data and points with error bars experimental values.

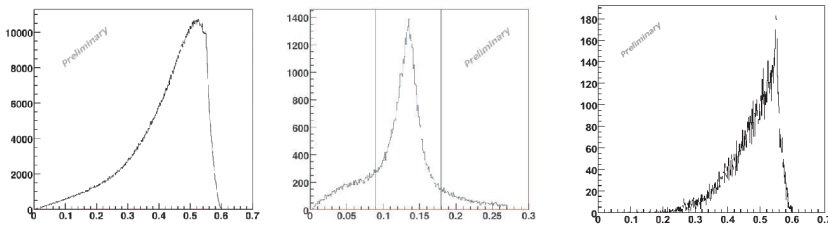


Fig. 3. **(left)** Missing mass of the  $pp \rightarrow ppX$  reaction. **(middle)** Invariant mass of two  $\gamma$  with cut lines. **(right)** Missing mass of two protons with the requirement of two  $\gamma$  in coincidence.

and straw tube trackers (FPC and MDC). Two protons were used to tag the  $\eta$  meson in the missing mass plot showed in Fig. 3 left (here we present data only from one run). The two  $\gamma$  originating from the  $\pi^0$  meson decay were registered in the electromagnetic calorimeter. The invariant mass of these  $\gamma$  is required to be close to the mass of the  $\pi^0$  (Fig. 3 middle). Requiring two  $\gamma$  in coincidence with the two protons gives the missing mass as shown in Fig. 3 right.

#### 4. Outlook

The work for both  $pd$  and  $pp$  data will be continued in order to obtain two independent determinations of the Dalitz plot density for the  $\eta \rightarrow \pi^+\pi^-\pi^0$ . This includes estimating systematical errors as well as tuning Monte Carlo simulation.

#### References

1. C. Roiesnel, T. Truong, *Nucl. Phys. B* **187**, 293 (1981).
2. J. Gasser, H. Leutwyler, *Nucl. Phys. B* **250**, 539 (1985).
3. R. Dashen, *Phys. Rev.* **183**, 1245 (1969).
4. F. Ambrosini *et al.* [KLOE Collaboration], *JHEP* **05**, 006 (2008).
5. WASA-at-COSY Collaboration, arXiv:nucl-ex/0411038, (2004).