# Testing CPT symmetry in ortho-positronium decays with the J-PET facility

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Total-body J-PET and Theranstic Center group photo, 2021-04-26 In front of Faculty of Physics, Astronomy and Computer Science of Jagiellonian University

# Motivation: discrete symmetry tests with o-Ps decays



## Testing discrete symmetries with angular correlations in o-Ps ${\rightarrow}3\gamma$ decays

Measurement the expectation value of the symmetry odd-operators



$$\begin{pmatrix} \hat{O} \\ \stackrel{?}{=} 0 & \text{for an odd operator} \\ \Leftrightarrow \mathcal{CPT}(\hat{O}) = -1 \\ \Leftrightarrow \mathcal{T}(\hat{O}) = -1 \end{cases}$$

 $|\vec{k}_1| > |\vec{k}_2| > |\vec{k}_3|$ 

4

### **Previous measurements**



 $C_{CPT} = (2.6 \pm 3.1) \times 10^{-3}$ 



Limiting positron emission direction 1 Mbq  $\beta^+$  emitter activity  $4\pi$  detector but low angular resolution

Yamazaki et al. PRL 104 (2010) 083401  $(\vec{S} \cdot \vec{k_1})(\vec{S} \cdot (\vec{k_1} \times \vec{k_2}))$ 

 $C_{CP} = (1.3 \pm 2.1 \pm 0.6) \times 10^{-3}$ 



Polarized o-Ps using external B field Inclusive measurement Only certain angular configurations

# Motivation: discrete symmetry tests with o-Ps decays

- Discrete symmetries are scarcely tested with leptonic systems:
  - Neutrino oscillations: Dirac phase, δCP ~3σ level [T2K, Nature 580 (2020) 339]
  - Electron EDM < 1.1x10<sup>-29</sup> [ACME, Nature 562 (2018) 355]



symmetries tests can be made with a very high precision limited, only by the effects due to the weak interaction: 10<sup>-14</sup> and photon-photon interaction: 10<sup>-9</sup>. (Standard Model Calculations)
 [Phys. Rev. A 37, 3189 (1988), Z. Phys. C 41, 143 (1988), M. S Sozzi "Discreet Symmetries and CP violation"]

### o-Ps production in J-PET with an annihilation chamber



[D. Kaminska, et a.l., Eur. Phys. J. C (2016) 76:445]

## o-Ps production in J-PET with an annihilation chamber



### J-PET detector at Jagiellonian University in Kraków, Poland



[P. Moskal et al., Acta Phys. Polon. B47 (2016) 509; G. Korcyl, et al., IEEE Trans. Med. Imag. 37, 2526 (2018)]

### o-Ps spin estimation:

\* e<sup>+</sup> spin estimated event-by-event recording multiple geometrical configurations
\* effective polarization depends

on o-Ps $\rightarrow$ 3 $\gamma$  vertex resolution \* vacuum in the chamber assures that e<sup>+</sup> is not going to decay before reaching inner wall

$$P_{\text{o-Ps}} = \frac{2}{3}P_{e+}$$

$$\beta^+$$

$$\beta^+$$

$$S$$

$$S$$

$$F_{e+} \approx \frac{v}{c} \cdot \frac{1}{2}(\cos \alpha + 1)$$

$$\mathbf{P}_{\rm e^+} = (N_{\rm +1/2}^{\rm e^+} - N_{\rm -1/2}^{\rm e^+}) / (N_{\rm +1/2}^{\rm e^+} + N_{\rm -1/2}^{\rm e^+})$$

### o-Ps spin determination and o-Ps $\rightarrow$ 3 $\gamma$ decays reconstruction in J-PET



The decay point  $(\mathbf{x',y'})$  in the decay plane and time **t** is an intersection of 3 circles, each corresponding to a possible origin points of the incident  $\gamma$ 

#### <u>o-Ps→3y decays reconstruction:</u>

\* Trilateration-based reconstruction to determine the o-Ps annihilation point





 $(T_i - t)^2 c^2 = (X'_i - x')^2 + (Y'_i - y')^2, \quad i = 1, 2, 3$ 

## **J-PET vs previous measurements**



 $C_{CPT} = (2.6 \pm 3.1) \times 10^{-3}$ 



 $\begin{array}{l} \mbox{Limiting positron emission direction} \\ 1 \mbox{ Mbq } \beta^{+} \mbox{ emitter activity} \\ 4 \pi \mbox{ detector but low angular resolution} \end{array}$ 



Recording multiple geometrical configurations

I-PET

e+ spin estimated event-by-event  $P_{e+} \approx \frac{v}{c} \cdot 0.91$  Yamazaki et al. PRL 104 (2010) 083401  $(\vec{S} \cdot \vec{k_1})(\vec{S} \cdot (\vec{k_1} \times \vec{k_2}))$ 

 $C_{CP} = (1.3 \pm 2.1 \pm 0.6) \times 10^{-3}$ 



Polarized o-Ps using external B field Inclusive measurement Only certain angular configurations

- Plastic scintillators = fast timing

   → using high β<sup>+</sup> emitter activity
   (tested up to 10 Mbq)
- Recording all 3 annihilation photons
- Angular resolution at 1° level

Courtesy of A. Gajos

### Identification of o-Ps $\rightarrow$ 3 $\gamma$ annihilation events in J-PET



[S. Sharma, eta al., EJNMMI Phys. 7, 39 (2020)]

Using total Time Over Threshold (TOT) of PMT signals from a scintillator strip which corresponds to y deposited energy





 $\leftarrow$  o-Ps presence in positron lifetime distribution

### **Background subtraction**

**Secondary Compton scatterings:** 

$$*\,\delta_{ij} = |d_{ij} - c\Delta t_{ij}|$$

computed for each pair of annihilation photon candidates i and j (i,j=1,2,3)



\* distance between the  $\beta$ + source location and the closest hypothetical  $2\gamma$  annihilation point on a LOR between two recorded photon interactions

 \* the sum of the two smallest angles between azimuthal coordinates of the recorded γ interaction points



### **Determination of the CPT - asymmetric observable**

$$O_{CPT} = \hat{S} \cdot \left(\vec{k}_1 \times \vec{k}_2\right) / \left|\vec{k}_1 \times \vec{k}_2\right| = \cos \phi$$

the angle between the direction of initial spin of the o-Ps atom and the normal to the decay plane

J-PET is sensitive to the full range of this operator





$$O_{CPT} = \hat{S} \cdot \left(\vec{k}_1 \times \vec{k}_2\right) / \left|\vec{k}_1 \times \vec{k}_2\right| = \cos \phi$$

the angle between the direction of initial spin of the o-Ps atom and the normal to the decay plane





### ARTICLE

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OPEN

# Testing CPT symmetry in ortho-positronium decays with positronium annihilation tomography

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### P. Moskal, et al., Nature Commun. 12, 5658 (2021)



Check for updates

## Summary



### **Summary and Perspectives**

- With J-PET scanner, we are able to perform exclusive measurement of ortho-positronium (o-Ps) annihilation into 3 photons
  - o-Ps spin event-by-event estimation
  - o-Ps→3γ decays reconstruction including determination of the annihilation point in an extensive-size medium

### Sub-permil precision of the CPT test reached with the first J-PET measurement (26 days): over factor of 3 better than the previous results

J-PET aims at the sensitivity of the CP and CPT symmetry tests at the level of 10<sup>-5</sup> with the pending improvements to the setup



 $C_{\rm CPT} = \langle O_{\rm CPT} \rangle / P = 0.00067 \pm 0.00095$ 

### **Summary and Perspectives**







new design of the annihilation chamber with spherical geometry, increasing the o-Ps formation probability by a factor of  $\sim$ 1.5

## **Summary and Perspectives**





additional densely packed layer of plastic scintillators with a fully digital readout -> increase of detection efficiency by factor of 64



# The first positronium imaging of a phantom built from cardiac myxoma and adipose tissue P. Moskal, et al., Science Advances 2021; 7 : eabh4394



# Thank you for your attention











# **BACKUP SLIDES**





# 192 module prototype





85 cm diameter,50 cm FOV4 constant thresholddiscrimination todetermine time andenergy of interaction



# 24 module prototype





time and angular resolutions of 100 ps and 0.4

- 74 cm diameter
- 50 cm FOV
- 4 SiPM per scintillator side
- 2 constant threshold per SiPM
- modular design 312 strips, 24 modules
- digital data at the module output
- very light ~60 kg





# **Discrete symmetry tests in positronium decays**

Measurement the expectation value of the symmetry odd-operators



example



# **Charge conjugation (C)**

The best limit in systems of quarks

STREET, STREET

$$\pi^0 \to 3\gamma$$
 is forbidden  $\frac{\pi^0 \to 3\gamma}{\pi^0 \to 2\gamma} < 3.1 \times 10^{-8}$  90% cl

CPT symmetry implies the equality of the masses, widths, etc. of a particle and its antiparticle

$$2\frac{|m_{K^0} - m_{\overline{K}^0}|}{(m_{K^0} + m_{\overline{K}^0})} < 6 \times 10^{-19}, \qquad 2\frac{|\Gamma_{K^0} - \Gamma_{\overline{K}^0}|}{(\Gamma_{K^0} + \Gamma_{\overline{K}^0})}$$

 $= (8 \pm 8) \times 10^{-18}$ 

- CP violation is equivalent to T violation
- CP violation asymmetry between matter and antimatter in our universe
- CP violation mechanism is introduced by the quark mixing described by the complex Cabibbo - Kobayashi - Maskawa matrix with none nonzero phase
- CP violation observed first for neutral kaons

CP violation J. H. Christenson et al., Phys. Rev. Lett. 13, 138 (1964).

$$|K_{\rm s}\rangle = \frac{1}{\sqrt{2}} \left[ |K^0\rangle + |\overline{K}^0\rangle \right] \text{ with } \mathcal{CP} = 1$$
$$|K_{\rm L}\rangle = \frac{1}{\sqrt{2}} \left[ |K^0\rangle - |\overline{K}^0\rangle \right] \text{ with } \mathcal{CP} = -1$$
$$\mathsf{BR} (\mathsf{K}_{\rm L} - \mathsf{R}^+ \pi^- / \mathsf{K}_{\rm L} - \mathsf{all}) \approx 2 \cdot 10^{-3}$$
$$|K_L\rangle = \frac{1}{\sqrt{1 + |\epsilon|^2}} \left( |K_2\rangle + \epsilon |K_1\rangle \right)$$
$$|K_S\rangle = \frac{1}{\sqrt{1 + |\epsilon|^2}} \left( |K_1\rangle - \epsilon |K_2\rangle \right)$$

