

DIRECT CP-VIOLATION AND RARE K DECAY PERSPECTIVES | Augusto Ceccucci / CERN

SETTING THE SCENE

Quark mixing and CP-Violation has been a very active area of investigation over the past decades

Owing to the last round of experiments in K and B mesons, our understanding is now completely compatible with the existence of "just" one phase in the Cabibbo-Kobayashi-Maskawa mixing matrix

The precision of the tests in the quark sector is improving thanks to the interplay of theory and experiments

Flavour transitions are so sensitive to short distance mechanisms that we need to press for quantitative tests of the Standard Model (SM)

Since the "directly" accessible energy frontier is limited (LHC for now) it is important to try explore the "zepto-universe", $O(10^{-21}$ m) "indirectly"

Baryon Asymmetry of the Universe (BAU)

Sakharov Conditions for BAU

Scanned at the American Institute of Physics

Andrei Sakharov (1967)

To allow the development of an asymmetry between matter and anti-matter

- **1. Violation of Baryonic Number**
- **2. Thermodynamic Non-equilibrium**

3. Violation of C & CP

Types of CP-Violation

$$
\left| M_{L} \right\rangle \propto p \left| M^{\circ} \right\rangle + q \left| \overline{M}^{\circ} \right\rangle
$$
\n
$$
\Delta F = 2 \qquad\n\left| A_{f} = \langle f | H | M \rangle, \quad \overline{A}_{f} = \langle f | H | \overline{M} \rangle \right\rangle
$$
\n
$$
\Delta F = 1
$$
\n
$$
\left| M_{H} \right\rangle \propto p \left| M^{\circ} \right\rangle - q \left| \overline{M}^{\circ} \right\rangle
$$
\n
$$
\Delta F = 2 \qquad\n\left| A_{\overline{f}} = \langle \overline{f} | H | M \rangle, \quad \overline{A}_{\overline{f}} = \langle \overline{f} | H | \overline{M} \rangle \right\rangle
$$
\n
$$
\Delta F = 1
$$

1. CP Violation in mixing $|q/p| \neq 1$ (indirect) 2. CP Violation in decays $|\overline{A}_{\overline{f}}/A_f| \neq 1$ (direct) 3. CP Violation in the interference

CP Violation

V.L.Fitch R.Turlay J.W.Cronin J.H.Christenson *Phys. Rev. Lett. 13 (1964) 138*.

6 $\frac{0}{\sqrt{\pi}}$ $\frac{1}{\pi}$ $\frac{1}{\sqrt{V}}$ $\frac{0}{\sigma}$ $\pi^+\pi^ \Gamma(\overline{V}^0)$ $\sim (5.5 \pm 0.5) \times 10$ $(K^0 \to \pi^+ \pi^-) + \Gamma(\overline{K}^0 \to \pi^+ \pi^-)$ $(K^0 \to \pi^+ \pi^-)$ - $\Gamma(\overline{K}^0 \to \pi^+ \pi^-)$ (5.5+0.5) $\times 10^{-7}$ $+\pi$ ⁻ \sqrt{V} ⁰ π ⁺ π ⁻ $+\pi$ ⁻ $\Gamma(\overline{V}^0 \rightarrow \pi^+\pi^ \pm$ 0.5) \times $\Gamma(K^0 \to \pi^+ \pi^-) + \Gamma(\overline{K}{}^0 \to$ $\Gamma(K^0 \to \pi^+ \pi^-)$ - $\Gamma(\overline{K}{}^0 \to$ $\pi \pi \rightarrow + 1$ (K) $\rightarrow \pi \pi$ $\pi \pi \rightarrow -1$ (K) $\rightarrow \pi \pi$ $K^0 \to \pi^+ \pi^-) + \Gamma(\overline K)$ $K^0 \to \pi^+ \pi^-) - \Gamma(\overline K$

BaBar + Belle: $\sin 2 \beta = 0.672 + 0.023$

Mixing Decay Interference

Quark masses and mixing

 The masses and mixings of quarks have a common origin in the standard model (SM): they arise from the Yukawa interactions with the Higgs condensate

 $\mathcal{L}_Y = -Y_{ij}^d \overline{Q_{Li}^I} \phi \, d_{Rj}^I - Y_{ij}^u \overline{Q_{Li}^I} \epsilon \phi^* u_{Rj}^I + \text{h.c.}$

- \odot When ϕ acquires a VEV we get the masses of the quarks
- The diagonalization yields the physical states. As a result the charged currents couples to the physical quarks as:

$$
\frac{-g}{\sqrt{2}} \left(\overline{u_L}, \overline{c_L}, \overline{t_L} \right) \gamma^{\mu} W_{\mu}^+ V_{\text{CKM}} \begin{pmatrix} d_L \\ s_L \\ b_L \end{pmatrix} + \text{h.c.},
$$

 \circ $\sqrt{\frac{C_{KM}}{C_{KM}}}$ is a 3 x 3 complex matrix know as the Cabibbo, Kobayashi, Maskawa matrix

Cabibbo-Kobayashi-Maskawa (CKM) Quark Mixing If V is unitary:

$$
V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}
$$

PDG 2014

nxn real parameter 2n-1 unphysical pahses n(n-1)/2 rotation angles (n-1)(n-2)/2 complex phases

 $|V_{ud}| = 0.97425 \pm 0.00022$ 0⁺ \rightarrow 0⁺ super-allowed nuclear β decays $|V_{us}| = 0.2253 \pm 0.0008$ Kaon semi-leptonic and leptonic decays $|V_{\text{cd}}| = 0.225 \pm 0.008$ semi-leptonic D decays and neutrino/antineutrino
 $|V_{\text{cs}}|$ =0.986 ± 0.016 $\text{Average of semi-leptonic D and leptonic D}_{\text{s}}$ decay $|V_{cs}|$ =0.986 ± 0.016
 $|V_{cs}|$ = (41.1 ± 1.3) × 10⁻³ Combination of exclusive and inclusive B decays $|V_{cb}| = (41.1 \pm 1.3) \times 10^{-3}$ Combination of exclusive and inclusive B decays $|V_{ub}| = (4.13 \pm 0.49) \times 10^{-3}$ Comb. of exclusive and inclusive charmless B de $|V_{ub}| = (4.13 \pm 0.49) \times 10^{-3}$ Comb. of exclusive and inclusive charmless B decays*
 $|V_{tb}| = 1.021 \pm 0.032$ Single top-quark production cross-section Single top-quark production cross-section

 V_{td} & V_{ts} accessible from FCNC processes (loops)

*But tension inclusive and exclusive determinations

One (of the six) Unitarity Relations

$$
V_{td}V_{tb}^* + V_{cd}V_{cb}^* + V_{ud}V_{ub}^* = 0
$$

Constraints on the rho eta plane

PDG 2014

The unique measure of CP-Violation in the SM is the area of the Unitarity Triangle (Jarlskog invariant J)

$$
J = (2.96^{+0.20}_{-0.16}) \times 10^{-5}
$$

PROMISE OF NEW PHYSICS IN

NDS

Standard Model

"New Physics"

 $\mathcal{A}_{\mathsf{SM}}$

$$
A_{SM} + A_{NP} = K_{SM} \frac{\alpha_W}{4\pi} \frac{F_{CKM}}{M_W^2} + K_{NP} L \frac{F_{NP}}{\Lambda^2}
$$

- *L* is a possible loop factor
- \odot $K_{NP} \sim K_{SM}$
- *F_{NP}* is the NP Flavour coupling
- If *L*> $\alpha_{\rm W}/4\pi$ and $F_{\rm NP}$ > $F_{\rm SM}$ we can extract the NP scale Λ

Bounds on Λ in TeV ($c_{\rm NP} = 1$) Bounds on c_{NP} $(A = 1 \text{ TeV})$ Observables Operator Re Re Im Im 9.0×10^{-7} $(\bar{s}_L \gamma^{\mu} d_L)^2$ 9.8×10^{2} 1.6×10^{4} 3.4×10^{-9} Δm_K ; ϵ_K 2.6×10^{-11} 1.8×10^{4} 3.2×10^{5} 6.9×10^{-9} $(\bar{s}_R d_L)(\bar{s}_L d_R)$ $(\bar{c}_L \gamma^\mu u_L)^2$ 1.2×10^{3} 2.9×10^{3} 1.0×10^{-7} 5.6×10^{-7} Δm_D ; $|q/p|_D$, ϕ_D $(\bar{c}_R u_L)(\bar{c}_L u_R)$ 6.2×10^3 1.5×10^{4} 5.7×10^{-8} 1.1×10^{-8} $(\overline{b_L} \gamma^{\mu} d_L)^2$ 6.6×10^2 9.3×10^{2} 1.1×10^{-6} 2.3×10^{-6} $\Delta m_{B_{d}}$; sin(2 β) from $B_{d} \to \psi K$ $\frac{(\overline{b}_R d_L)(\overline{b}_L d_R)}{(\overline{b}_L \gamma^{\mu} s_L)^2}$ 2.5×10^3 3.9×10^{-7} 1.9×10^{-7} 3.6×10^3 1.4×10^{2} 2.5×10^{2} 1.7×10^{-5} 5.0×10^{-5} $\Delta m_{\mathcal{B}_{s}}$; $\sin(\phi_{s})$ from $B_{s} \to \psi \phi$ $(\bar{b}_R s_L)(\bar{b}_L s_R)$ 8.8×10^{-6} 4.8×10^{2} 8.3×10^2 2.9×10^{-6}

Isidori and Teubert, arXiv:1402.2844

B RARE DECAYS

- Few meson decays are particularly clean theoretically and so suppressed in the Standard Model that they provide a window to very short distance
- For B's I will just mention one example and the prospects for the next decades

- Observation of $B_s^0 \to \mu^+ \mu^-$ using combined CMS and LHCb dataset [arxiv:1411.4413], submitted to Nature
- $\mathcal{B}(B_s^0 \to \mu^+ \mu^-) = (2.79^{+0.66}_{-0.60}{}_{-0.19}{}^{\circ}) \times 10^{-9}$, 6.2 σ sign. (7.6 σ expected) $\mathcal{B}(B^0 \to \mu^+ \mu^-) = (3.94^{+1.58}_{-1.41}{}_{-0.24}) \times 10^{-10}$, 3.2σ sign. (0.8 σ expected)
- SM predictions [Bobeth et al., PRL 112 (2014) 101801] $\mathcal{B}(B_s^0 \to \mu^+ \mu^-) = (3.66 \pm 0.23) \times 10^{-9}$, compatible at 1.2σ $\mathcal{B}(B^0 \to \mu^+ \mu^-) = (1.06 \pm 0.09) \times 10^{-10}$, compatible at 2.2σ

KERKER E DAQ

ECFA STUDY ON HEAVY FLAVOUR

Expected precision on γ from tree decays

- 2025 appears to be the time for a good harvest
- A crossroad for B physics...

KAON RARE DECAYS

- K physics alone can fully constrain the $\mathcal{L}_{\mathcal{A}}$ CKM unitarity triangle.
- Comparison with B physics can provide \bullet description of NP flavour dynamics

K⁰ L[→] *0e e* $\mathcal{L}^0 \rightarrow \pi^0 \mu^+ \mu^-$

Study Direct CP-Violation

•**NA48/1 has measured the Indirect CP-Violating Contribution for both modes** •**S-L Constructive Interference preferred** •**CP-Conserving Contributions are negligible**

Direct CPV

K TeV: $K_L \rightarrow \pi^0$ ee

1999 data

PRL93, 021805 (2004)

• **One candidate in the signal box**

$$
BR(K_L \to \pi^0 \, ee \,) < 3.5 \times 10^{-10} \, @ 90\% CL
$$

• **Combining 1997 and 1999:**

Expected Background 0.99 ± 0.35 events

 K TeV: $K^o_L \rightarrow \pi^o \mu \mu$

 $\text{BR}(\text{K}^0_{\text{L}}\to\pi^0\mu\mu)$ $< 3.8\times 10^{10}$ (90% C.L.) [prl 86, 5425 (2001)]

 K^0 _{*S*} $\rightarrow \pi^0$ e⁺e \int and K^0 _S $\rightarrow \pi^0$ $\mu^+ \mu^-$

K^0 _{*L→π*⁰ee (μμ): **SM Branching Ratios**}

 $\frac{1}{10^{-4} \text{m}^{-1}}$ $(\times 10^{-12}$ **Thank to the NA48/1 measurements, the KL BR can now be predicted**

<u>Litatik to the NA40/Tilleasurements, the KL BK candidate $Br(\rm{K}_{L} \rightarrow \pi^0 \mu^+ \mu^-)$ $\sigma^0 \nu^+ \mu^-)$ </u>

Constructive

 1.1×10^{-11} $\sum_{e^+e^-}$ = 3.7^{+1.1} × 10 -11 $B_{a^+a^-}$ +1.1 $\times 10^{-11}$ $=3.7^{+1.1}_{-0.9}\times10$ $+0.3 \times 10^{-11}$

$$
B_{\mu^+\mu^-} = 1.5^{+0.3}_{-0.3}\times10^{-1}
$$

now favored by two independent analyses*

Destructive

 0.7×10^{-11} $\sum_{e^+e^-}$ = 1.7^{+0.7} × 10 0.2×10^{-11} $B_{\mu^+\mu^-} = 1.0^{+0.2}_{-0.2} \times 10^{-1}$ $B_{a^+a^-}$ $+0.7 \times 10^{-11}$ $=1.7^{+0.7}_{-0.6}\times10$ $+0.2 \times 10^{-11}$ $m_{\mu^+\mu^-} = 1.0^{+0.2}_{-0.2} \times 10^{-1}$

*G. Buchalla, G. D'Ambrosio, G. Isidori, Nucl.Phys.B672,387 (2003) *S. Friot, D. Greynat, E. de Rafael, hep-ph/0404136 *

CP violating: $\text{K}_\text{S}{\rightarrow}\pi^0\pi^0\pi^0$ never observed so far !

 $SM \Gamma(K_S \rightarrow 3\pi^0) = \Gamma(K_L \rightarrow 3\pi^0) |\eta_{000}|^2 \Rightarrow BR(K_S \rightarrow 3\pi^0) \sim 2 \times 10^{-9}$

KLOE/KLOE-2: Phys. Lett. B 723 (2013) 54

 $BR(K_S \to 3\pi^0) \leq 2.6 \times 10^{-8}$ at 90\% C.L.

Factor of five better than previous results

 $\epsilon + \epsilon'_{000} = |\eta_{000}| = \left| \frac{A(K_S \to 3\pi^0)}{A(K_L \to 3\pi^0)} \right| = \sqrt{\frac{\tau_L}{\tau_S} \frac{BR(K_S \to 3\pi^0)}{BR(K_L \to 3\pi^0)}} \leq 0.0088$ at 90% C.L.

Courtesy P. Moskal

KLOE-2 has a chance to observe K_S --> $\pi^0 \pi^0 \pi^0$ decay **for the first time in the near future**

 $Br(K_L^0 \to \pi^0 \nu \overline{\nu})$

HOLY GRAIL OF FLAVOUR PHYSICS?

- Why it is so special:
- 1. Apart from a small admixture $(\epsilon_{\text{K}}$ 2.228 10⁻³), K^0 *L* is a CP eigenstate. Neglecting the CP-even state we can write:

$$
\langle \pi^0 \nu \overline{\nu} | A | K^0 \rangle \sim V_{td} V_{ts}^* X(x_t) + P_c(X) V_{cd} V_{cs}^*
$$

$$
\langle \pi^0 \nu \overline{\nu} | A | \overline{K}^0 \rangle \sim V_{td}^* V_{ts} X(x_t) + P_c(X) V_{cd}^* V_{cs}
$$

$$
|K_L^0\rangle \sim \frac{K^0 - \overline{K}^0}{\sqrt{2}}
$$

2. In taking the difference, the charm part (which is almost real) drops off and only the imaginary part of the top contribution remains!

> ${}^{0}\nu\overline{\nu}$ | A | K_{L}^{0} > ~ Im $V_{td}V_{ts}^{*}X(x_{t})$ $<$ $\pi^{0}\nu\overline{\nu}$ | A | K_{L}^{0} >~ ${\rm Im} V_{td}V_{ts}^{*}X$ (x_{t}

- 3. The main experimental background $(K^0_L \to \pi^0 \pi^0)$ is suppressed by CP conservation !
- 4. The very long life time of the K^0 _L makes the interesting partial width "measurable" $(Br-O(10^{-11}))$

 $Br(K_L^0 \rightarrow \pi^0 \nu \overline{\nu})$

Formulas from A.J. Buras et al. RMP 80, 2008

$$
Br(K_L^0 \to \pi^0 \nu \overline{\nu})
$$

Formulas from A.J. Buras et al. RMP 80, 20

$$
Br(K_L^0 \to \pi^0 \nu \overline{\nu}) = \kappa_L \times \left(\frac{\text{Im }\lambda_t}{\lambda^5} X(x_t)\right)^2
$$

$$
\kappa_L = (2.231 \pm 0.013) \times 10^{-10} \left[\frac{\lambda}{0.225}\right]^8
$$

Numerical example:
 $Im V_{ud} V_{us}^* = sin \beta_K |V_{ud} V_{us}^*| \sim 1$
$$
X(x_t) \sim 1.44
$$

$$
\frac{|V_{us}|}{|V_{ud}|^2 + |V_{us}|^2}
$$

$$
Br(K_L^0 \to \pi^0 \nu \overline{\nu}) \sim 2.3
$$

EXPERIMENT: BR<2.6 10⁻⁸ 90%CL (E391a - 1
NEXT EXPERIMENT: KOTO (E14, J-PARC)

$$
\lambda_{i} = V_{id}V_{is}^{*} \qquad \qquad \text{Im}V_{td}V_{ts}^{*} = \sin \beta_{K} |V_{td}V_{ts}^{*}| \sim 1.29 \times 10^{-4}
$$
\n
$$
X(x_{t}) \sim 1.44
$$
\n
$$
\lambda = \frac{|V_{us}|}{\sqrt{|V_{cs}|^{2} + |V_{cs}|^{2}}} \qquad \qquad \text{Br}(K_{L}^{0} \to \pi^{0}V\overline{V}) \sim 2.3 \times 10^{-11}
$$

 $|V_{ud}|^2 + |V_{us}|^2$

 $|V_{ud}|^2 + |V|$

 $_{ud}$ | $+$ | V_{us}

EXPERIMENT: BR<2.6 10-8 90%CL (E391a - KEK)

Pure CsI recovered from FNAL KTeV Experiment

Current SES based on 100 h run in 2013 (Preliminary): 1.29×10^{-8}

Expect "nominal" beam intensity in 2017

$$
Br(K^+\to\pi^+\nu\overline{\nu}) = \kappa_+(1+\Delta_{EM})\times
$$

$$
\left[\left(\frac{\text{Im }\lambda_t}{\lambda^5}X(x_t)\right)^2 + \left(\frac{\text{Re }\lambda_c}{\lambda}P_c(X) + \frac{\text{Re }\lambda_t}{\lambda^5}\right)^2\right]
$$

$$
\kappa_+ = (5.173 \pm 0.025) \times 10^{-11} \left[\frac{\lambda}{0.225}\right]^8
$$

$$
\lambda = \frac{|V_{us}|}{\sqrt{|V_{ud}|^2 + |V_{us}|^2}}
$$
\n
$$
\kappa_{+} = r_{K^{+}} \cdot \frac{3\alpha^2 Br(K^{+} \to \pi^0 e^{+} \nu)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8
$$
\n
$$
\lambda_{i} = V_{id} V_{is}^{*}
$$

Formulas from A.J. Buras et al. RMP 80, 2008

$Br(K^+\to\pi^+\nu\overline{\nu})$

$$
Br(K^{+} \to \pi^{+}V\overline{V}) \propto 1.56 \times 10^{-4} \times
$$
\n
$$
\left[V_{td}V_{ts}^{*} \right]^{2} X(x_{t})^{2} + 2\lambda^{5} P_{c}(X) \left| V_{td}V_{ts}^{*} \right| X(x_{t}) \cos \beta_{K} + \lambda^{10} P_{c}(X)^{2} \right] \approx
$$
\n
$$
\left[4.40 + 3.68 + 0.87 \right] \times 10^{-11} =
$$
\n
$$
8.95 \times 10^{-11}
$$

The charm- top-quark interference term is comparatively large

$$
\cos \beta_K = \cos \beta - \beta_s \approx 0.94
$$

For this set of values the m_c the parametric uncertainty is: δ Br/Br ~ 0.68 δ P_c/P_c

 $|V_{td}V_{ts}^*|$ ~ 3.69×10⁻⁴ (PDG 2014)

 $X(x_t) \sim 1.44$ (Buras et al.)

 $P_c(X) = 0.41 \pm 0.05$ (Buras et al.)

Kaon Rare Decays and NP

(courtesy by Christopher Smith)

Rare *K* **decay sensitivity to flavor violating** *Z'*

Sensitivity beyond direct searches

CHARGED K BEAMS

"Stopped"

- Work in Kaon frame
- High Kaon purity (Electro-Magneto-static Separators)
- **Compact Detectors**

"In-Flight"

- **-** Decays in vacuum (no scattering, no interactions)
- **RF** separated or Unseparated beams
- **Extended decay regions**

STATE OF THE ART: E787/E949 DECAYS AT REST

B(K^+ $\rightarrow \pi^+$ vv) = (1.73^{+1.15}_{-1.05}) x 10⁻¹⁰

PRL101, arXiv:0808.2459, AGS-E787/E949

 $\mathbf P$ π

 \mathbf{v}

 P_{K}

 $\Theta_{\pi\mathrm{K}}$

NA62 IN-FLIGHT TECHNIQUE

- Calorimetry to veto extra particles
- Very light trackers to reconstruct the *K +* and the $\pi^{\! *}$ momenta
- Full particle identification

CERN ACCELERATORS

$K^+ \to \pi^+ \nu \bar{\nu}$ Analysis Sensitivity (MC)

<u> TÉRN</u>

CERN NA62 EXPERIMENT

- Picture taken just before beam in 2014
- Beam time 2015: June 22 November 15

View of ECN3

Straw3 - LAV10 - MNP33 - Straw2 - LAV9

RICH Straw 4 and LAV11

GIGATRACKER (GTK) NA62 &

CERN (PH-DT, PH-ESE, PH-SME, EN,…) Ferrara, Louvain-la-Neuve, Torino

GIGATRACKER PERFORMANC NAG2

K12 Beam; Illumination of one GTK chip

NA62 STRAW TRACKER

CERN (PH-DT, PH-ESE, PH-SME) - JINR

RICH DATA

RICH-KTAG

- Events with only 1 track in the spectrometer reconstructed (within 40 ns)
- 10² muon rejection at trigger level \bullet

- Matching between track and RICH ring to study the particle content \bullet
- Positrons suppressed by the trigger \bullet

45

- Kaon decay modes reconstructed with the liquid Krypton calorimeter only (from $\mathbf x$ minimum bias data)
- Useful to measure the kinematic suppression factor, particle ID efficiency ... \mathbf{x}

Further NA62 K Physics Program

10/03/2015

IÉRN

SUMMARY

Fascinating adventure in between quark mixing, CP-Violation…

Number of fermion generations…

- Quarks and Leptons…
- High energy frontier…

Address fundamental question: How does the world work?

Rare *K* decays research offer a complementary way to see beyond SM with respect to high energy colliders

Ready to harvest un unprecedented amount of kaon decays…many questions are awaiting answers…

SPARE

Apply KTAG for kaon identification \bullet

KTAG candidate No KTAG candidate

KRYPTON READ OUT

Specifications/Tender : CERN PH-ESE,PH-SME Manufacturer: CAEN (ITALY) 14 bit FADC, 40 Ms, 32 ch / module **432 modules, 28 VME crates**

NA62 LAV: Frascati, Naples, Pisa & Rome 1 Collaboration

UNITARITY TRIANGLE FOR KAONS

 When the bd UT is used, the variables extracted from kaons are affected by an apparent parametric uncertainty

- The six UTs are all born equal (in the SM they have the same measure of CP-viola $\operatorname{tion},$ the Jarlskog invariant J $_{\mathsf{CP}}$)
- A remarkable feature is that in the *ds* UT

J_{CP}= 5.6 * sqrt(BR(K_L→ π^0 \vee ν))

This is a determination which is basically free from theoretical error (down to 1-2%)

It is to be compared with the current J_{CP} determination from the bd UT fit where the error ranges from 3% to 7% depending on the treatment of the errors

LARGE ANGLE VETOES

12 Electro-magnetic calorimeters A1-A11 in Vacuum (including PMTs) A12 in air Lead Glass counters from the LEP-OPAL ECAL (~2500 blocks)

All stations Installed

A Standard Model view of CP-Violation in Kaons

 \odot Neutral Kaon Mixing ($\pi\pi$, semi-leptonic)

$$
|\varepsilon| = \frac{G_F^2 f_K^2 m_K m_W^2}{12\sqrt{2}\pi^2 \Delta m_K} \hat{B}_K \{ \eta_1 S(x_c) \text{Im}(V_{cs} V_{cd}^*)^2 + \eta_2 S(x_t) \text{Im}(V_{ts} V_{td}^*)^2 + 2\eta_3 S(x_c, x_t) \text{Im}(V_{cs} V_{cd}^*) \}
$$

 $\epsilon = (2.228 \pm 0.011) \times 10^{-3}$

 \odot Neutral Kaon Decays into $\pi\pi$

$$
\left| \text{Re} \frac{\varepsilon'}{\varepsilon} \propto \text{Im}(V_{td} V_{ts}^*) \left[P^{1/2} - P^{3/2} \right] e^{i(\phi_{\varepsilon'} - \phi_{\varepsilon})} \right|
$$

PDG Average

$$
\mathop{\rm Re}\nolimits\frac{\varepsilon'}{\varepsilon} = (1.68 \pm 0.14) \times 10^{-3}
$$

Hierarchical Structure in powers of

$$
V_{CKM} = \begin{pmatrix} 1 - \lambda^2 / 2 & \lambda & A\lambda^3 (\rho - i\eta) \\ -\lambda & 1 - \lambda^2 / 2 & A\lambda^2 \\ A\lambda^3 (1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}
$$

 λ = sine of the Cabibbo angle

Wolfenstein parameterization

From Global fit (PDG review, 2014):

Imposing SM constraint (3 generation unitarity):

 $λ = 0.22535 ± 0.00065$ $\overline{A} = 0.811 + 0.022$ _{-0.012}

 ρ = 0.131 ^{+0.026} _{-0.013} η = 0.345 ^{+0.013} -_{0.014}

Some Constraints on the Unitarity Triangle

