Status of determining meson-nucleus potentials and the search for mesic states

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Outline:

- introduction: meson-nucleon interactions
- methods for determining meson-nucleus potentials
- potential parameters for K^+,K^0,K^-, η, η’ω,Φ - A interaction
- search for meson-nucleus bound states
- summary & outlook

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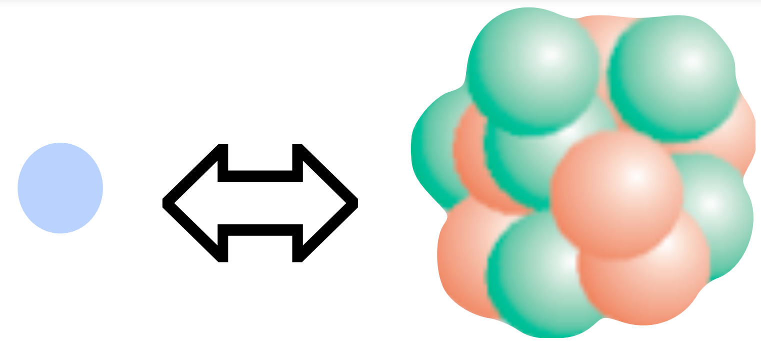
Cracow, Poland, June 5-9, 2017

HIC for FAIR
Helmholtz International Center

meson-nucleus interaction

meson  nucleus

1.) interaction attractive or repulsive ??

2.) if attractive, interaction strong enough
to form meson-nucleus bound state
exclusively bound by the strong interaction??

exotic nuclear configurations:
<u>nuclear physics:</u> states with excitation energies of several 100 MeV
<u>hadron physics:</u> investigate in-medium properties of mesons

mesons investigated: K^+,K^0,K^-,η,η',ω,Φ

meson-nucleus potential

H. Nagahiro, S. Hirenzaki, PRL 94 (2005) 232503

$$U(r) = V(r) + i\,W(r)$$

attractive ?
repulsive ?

absorption

$$V(r) = \Delta m(\rho_0) \cdot \rho(r)/\rho_0$$

$$W(r) = -\Gamma_0/2 \cdot \rho(r)/\rho_0$$
$$= -1/2 \cdot \hbar c \cdot \rho(r) \cdot \sigma_{inel} \cdot \beta$$

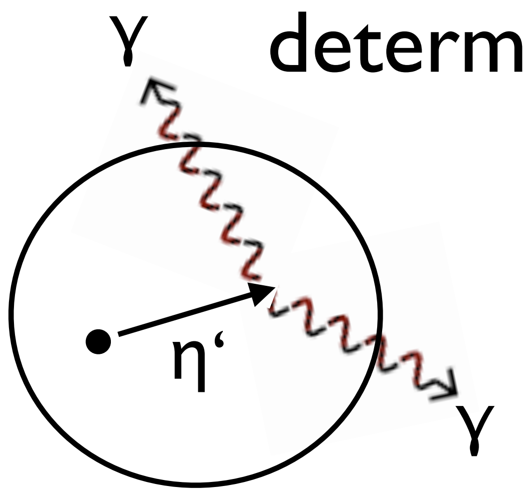
- line shape analysis
- excitation function
- momentum distribution
- meson-nucleus bound states

- transparency ratio measurement

$$T_A = \frac{\sigma_{\gamma A \rightarrow \eta' X}}{A \cdot \sigma_{\gamma N \rightarrow \eta' X}}$$

D. Cabrera et al., NPA733 (2004)130

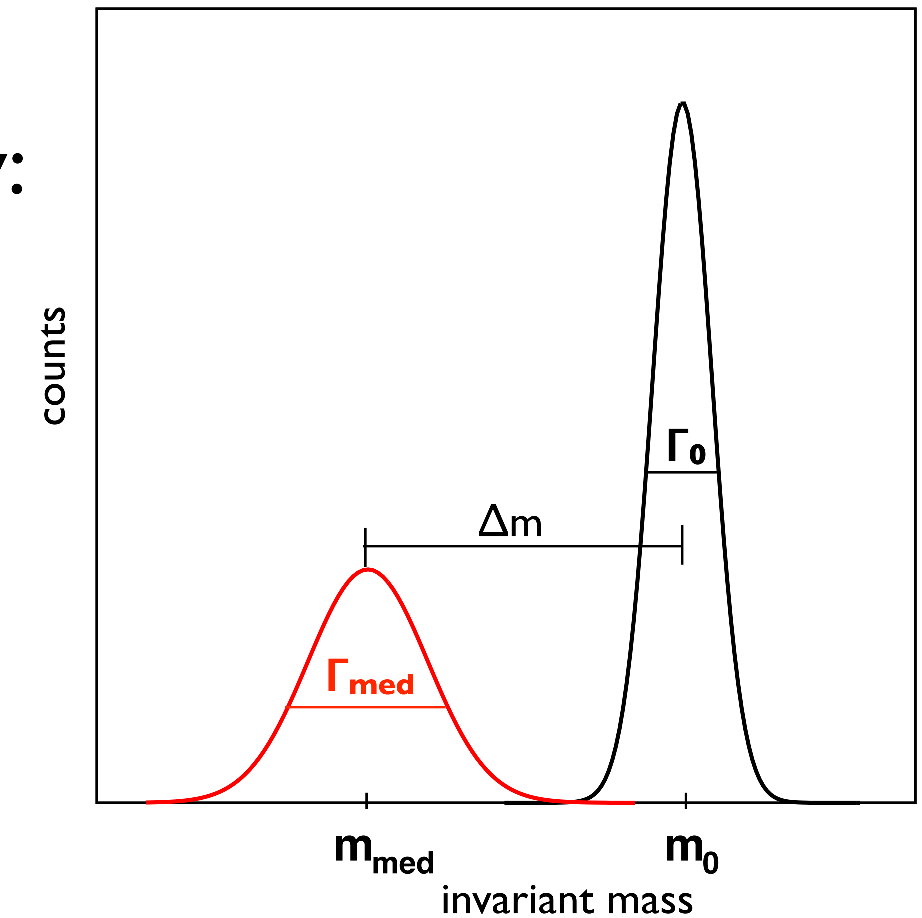
line shape analysis??



determine mass from in-medium decay:

e.g., $\eta' \to \gamma\gamma$

$$m = \sqrt{(p_1+p_2)^2}$$



probability for decay:

$$\frac{dP_{decay}}{dl} = \frac{mc}{p} \cdot \frac{1}{\hbar c} \cdot \Gamma_{decay} = 2.1 \cdot 10^{-5} /\mathrm{fm}$$

$$\Gamma_{\eta' \to \gamma\gamma} = 4.1 \cdot 10^{-3}\ \mathrm{MeV}$$

probability for absorption:

$$\frac{dP_{abs}}{dl} = \sigma_{abs} \cdot \rho(r) = 0.21/\mathrm{fm} \text{ at } \rho=\rho_0$$

$$\sigma_{abs} = 13\ \mathrm{mb}$$

$$\frac{P_{decay}}{P_{abs}} = 10^{-4}$$

10 000 times more likely
to get absorbed than to decay

more favourable decay/absorption ratio only at lower densities near the surface where in-medium modifications are reduced

sensitive to nuclear density at decay point

Determining the real part of the ω-nucleus potential: GiBUU transport model simulations

J.Weil, U. Mosel and V. Metag, PLB 723 (2013) 120 $\omega \rightarrow \pi^0\gamma$

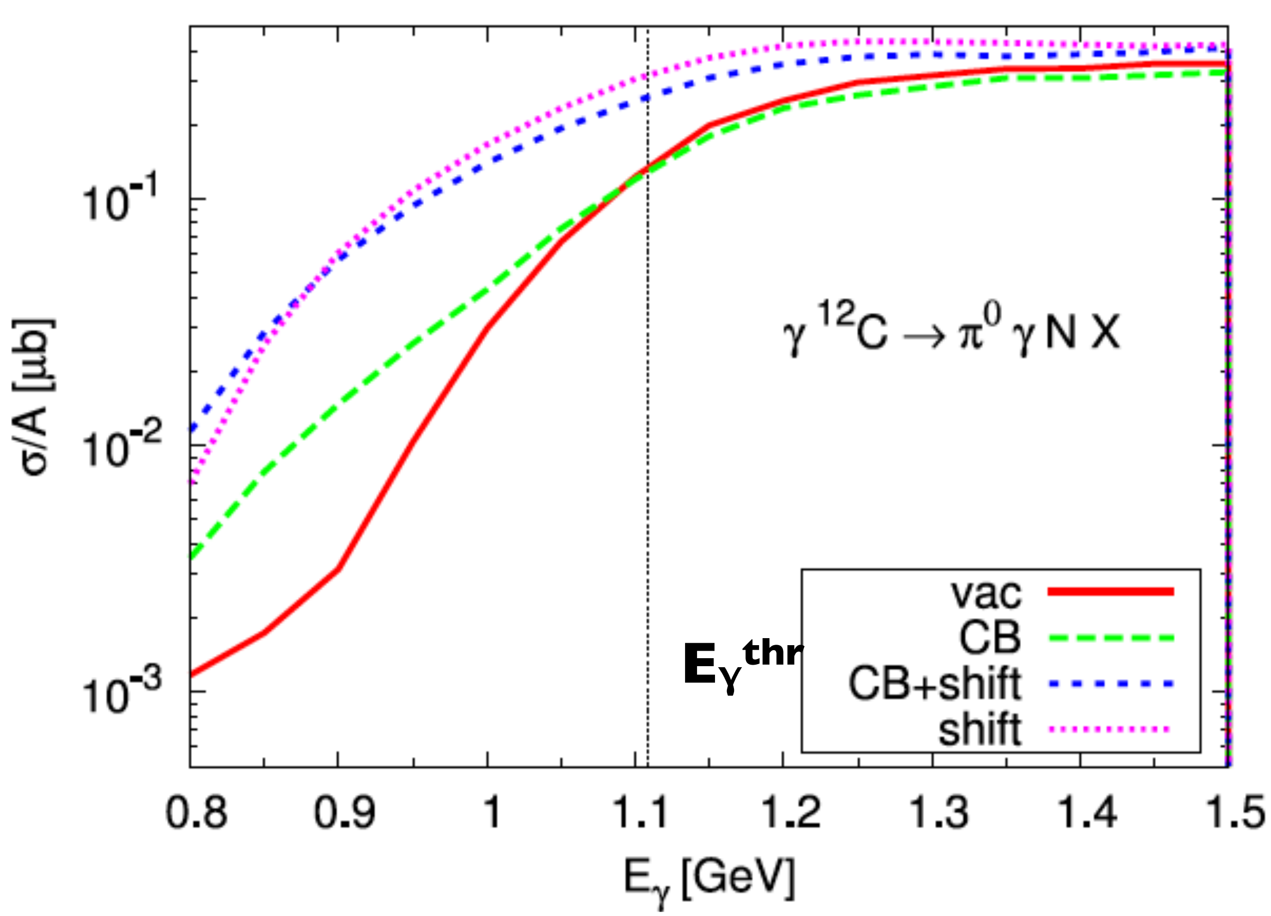
sensitive to nuclear density at production point

- measurement of the excitation function of the meson near threshold

in case of dropping mass -
higher meson yield for given $\sqrt{s}$
because of increased phase space
due to lowering of the production threshold

⇨ cross section enhancement

$\pi^0\gamma$ excitation function

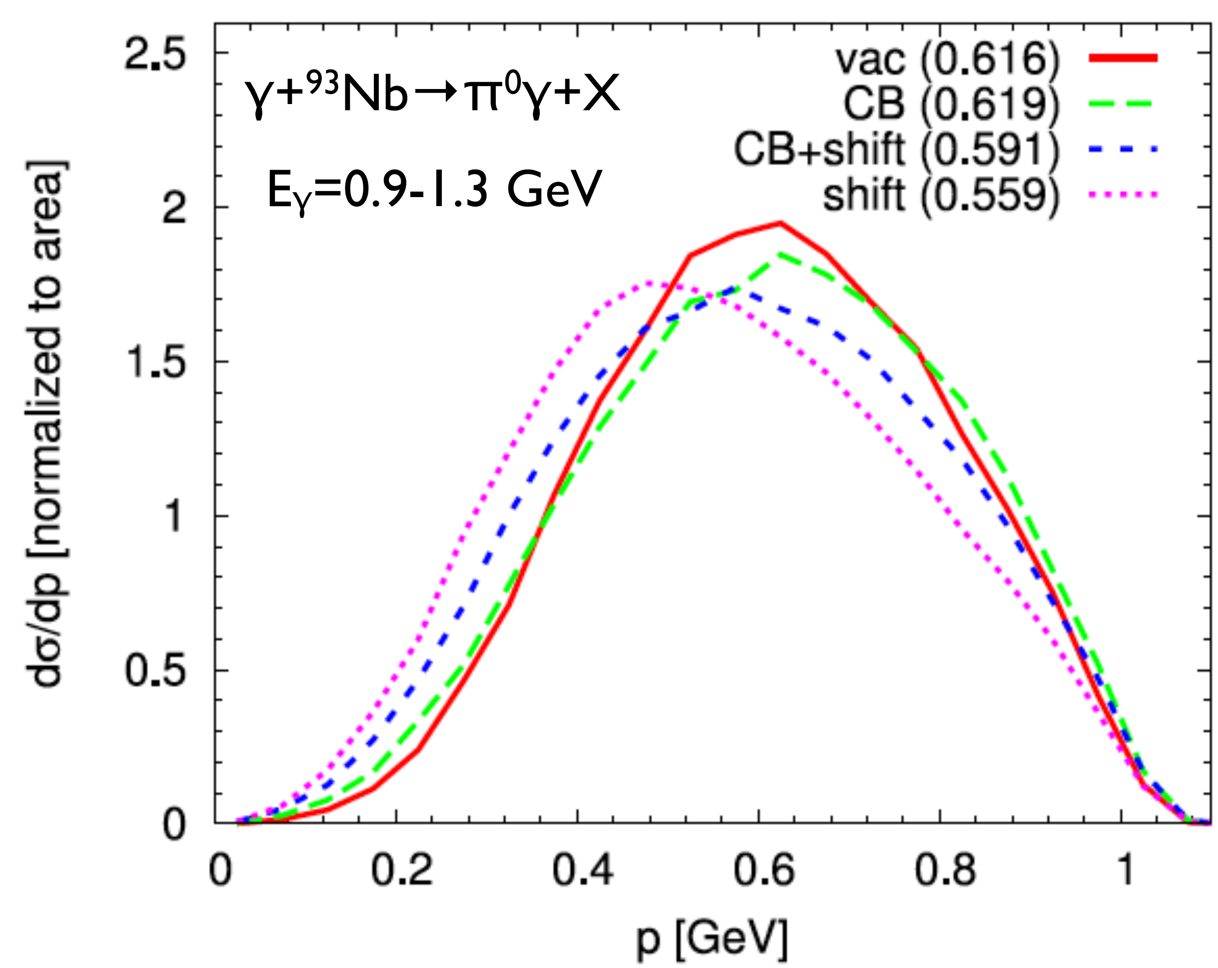


- momentum distribution of the meson:

in case of dropping mass - when leaving the nucleus hadron has to become on-shell; mass generated at the expense of kinetic energy

⇨ downward shift of momentum distribution

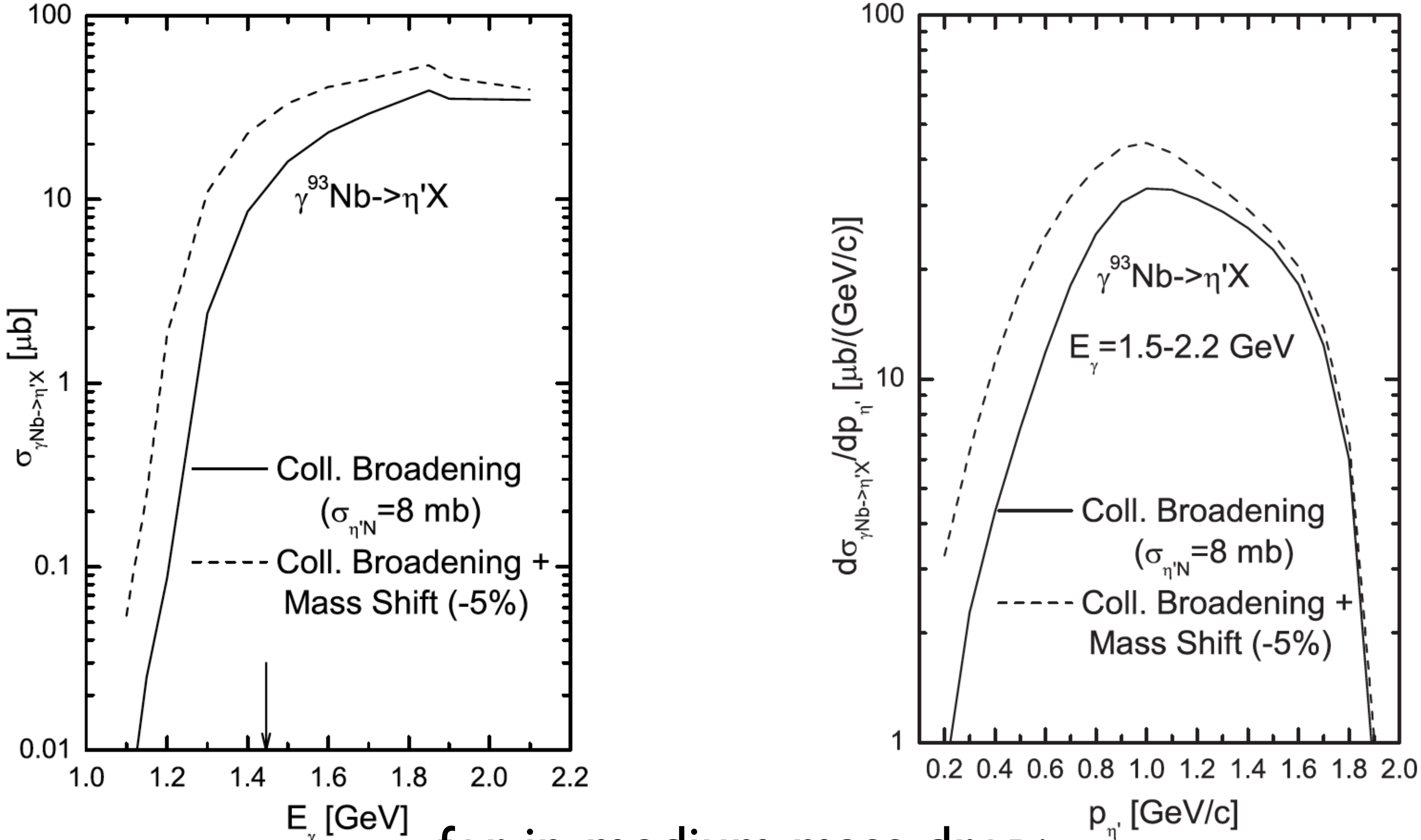
$\pi^0\gamma$ momentum distribution



Determining the real part of the η'-nucleus potential: collision model calculations

E. Ya. Paryev, J. Phys.G 40 (2013)025201

excitation function　　　　momentum distribution



for in-medium mass drop:

enhancement at low energies　　　　enhancement at low momenta

$$V(\rho=\rho_0) = -\Delta m\ (\rho=\rho_0)$$

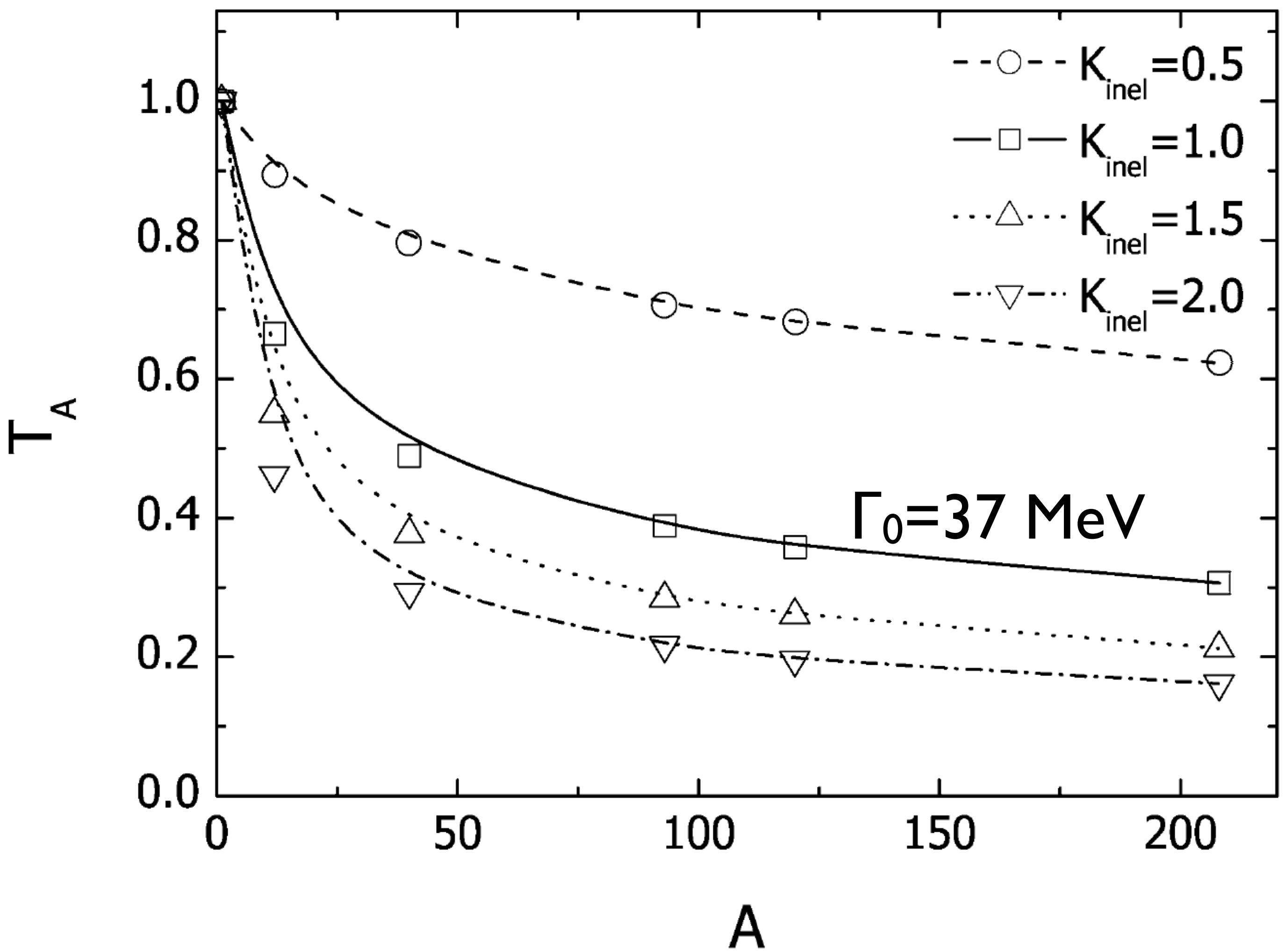
Determining the imaginary part of the meson-nucleus potential from transparency ratio measurements

$$T_A = \frac{\sigma_{\gamma A \to \eta' X}}{A \cdot \sigma_{\gamma N \to \eta' X}}$$

transport model calculation: GiBUU

P. Mühlich and U. Mosel, NPA 773 (2006) 156

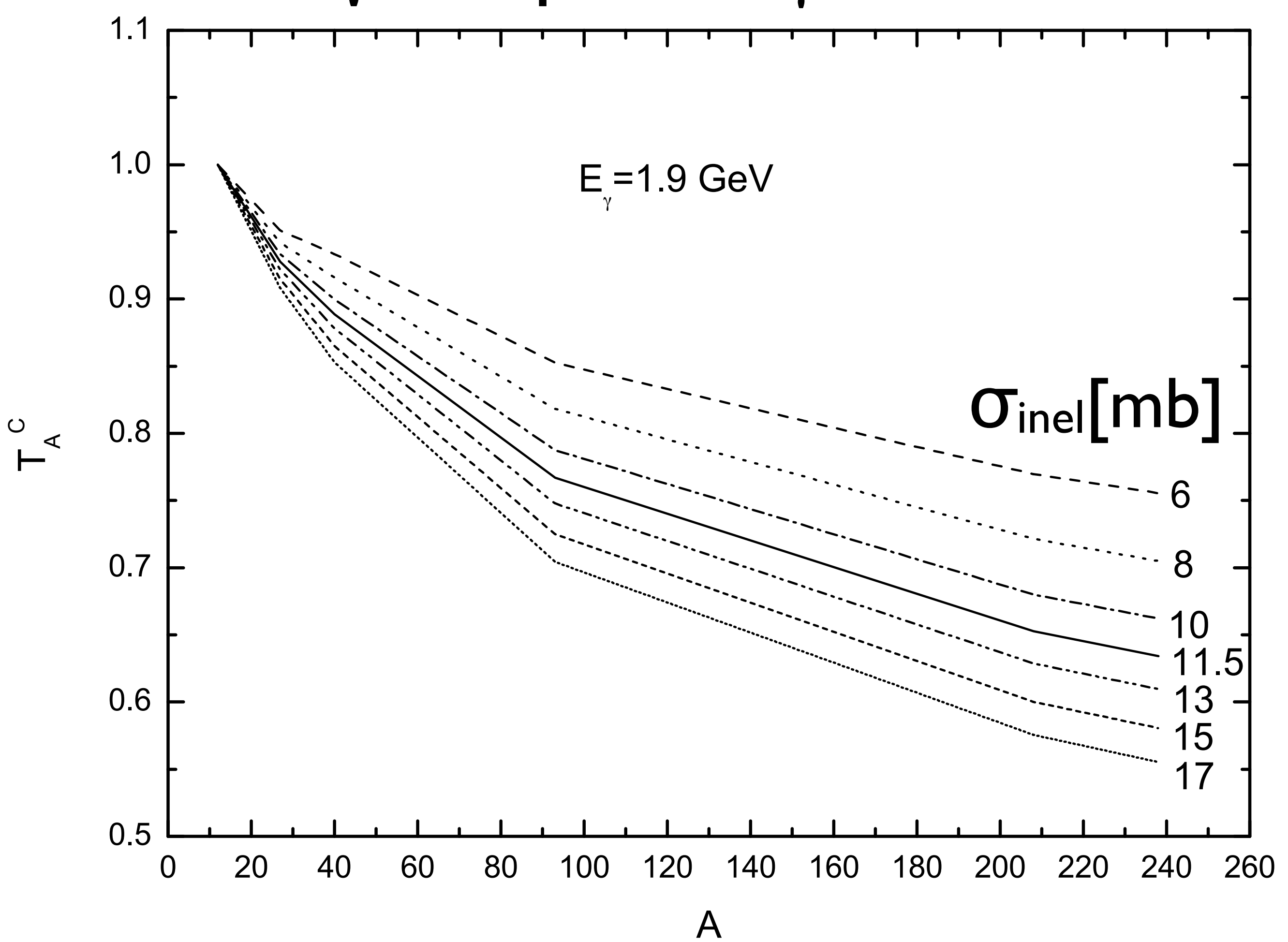
$\gamma A \to \omega X$ at E_γ=1.5 GeV



collision model calculation

E. Ya. Paryev, J. Phys.G 40 (2013)025201

$\gamma A \to \eta' X$ at E_γ=1.9 GeV



$$W(\rho=\rho_0) = -\Gamma/2\ (\rho=\rho_0) = -1/2 \cdot \hbar c \cdot \rho_0 \cdot \sigma_{inel} \cdot \beta$$

strategy for determining potential parameters

real part of meson-nucleus potential

measure meson excitation functions and/or momentum distributions

compare with transport and or collision model calculations
for different sets of V_0

$$\rightarrow V_0 = V(\rho=\rho_0)$$

imaginary part of meson-nucleus potential

measure transparency ratio $T_A(A,p)$

compare with transport and or collision model calculations
for different sets of Γ_{med}, σ_{inel}

$$\rightarrow \Gamma_{med}, \sigma_{inel} \rightarrow W_0 = W(\rho=\rho_{0;}\, p=0)$$

$$U(\rho=\rho_0) = V_0 + i\, W_0$$

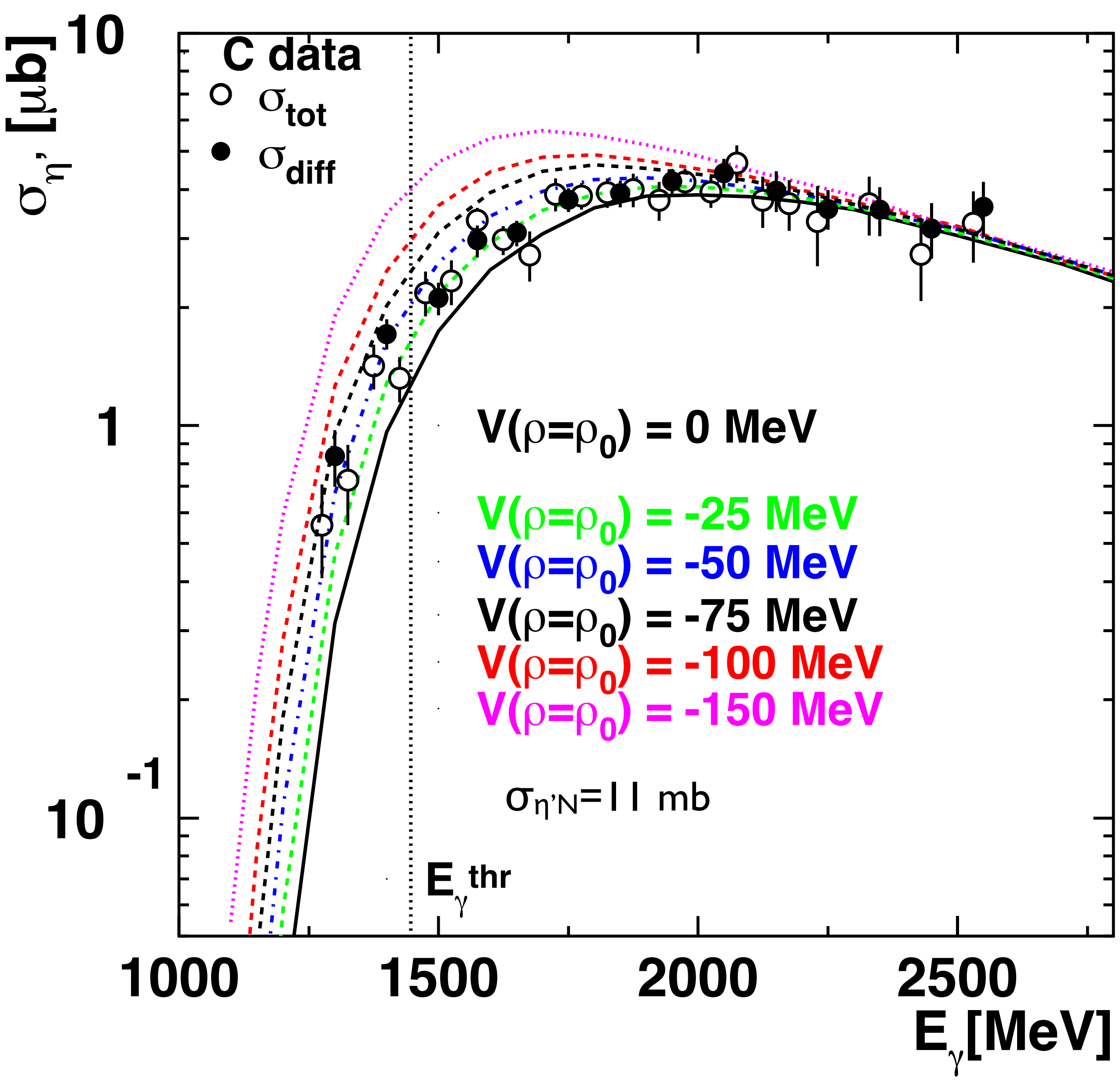
excitation function and momentum distribution for η' photoproduction off C

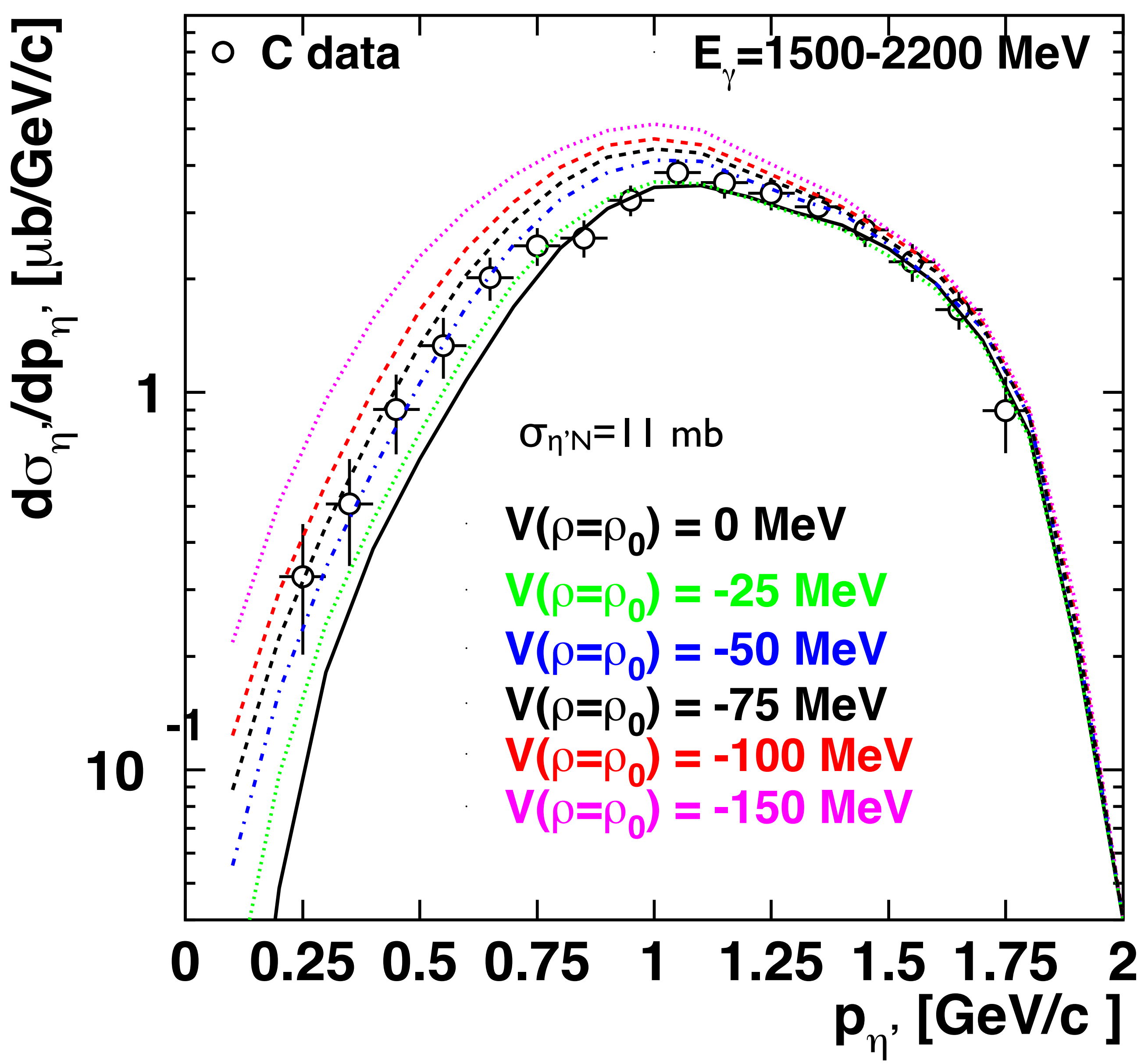
CBELSA/TAPS @ ELSA

$\gamma\, C \rightarrow \eta' X$

data: M. Nanova et al., PLB 727 (2013) 417

calc.: E. Paryev, J. Phys. G 40 (2013) 025201





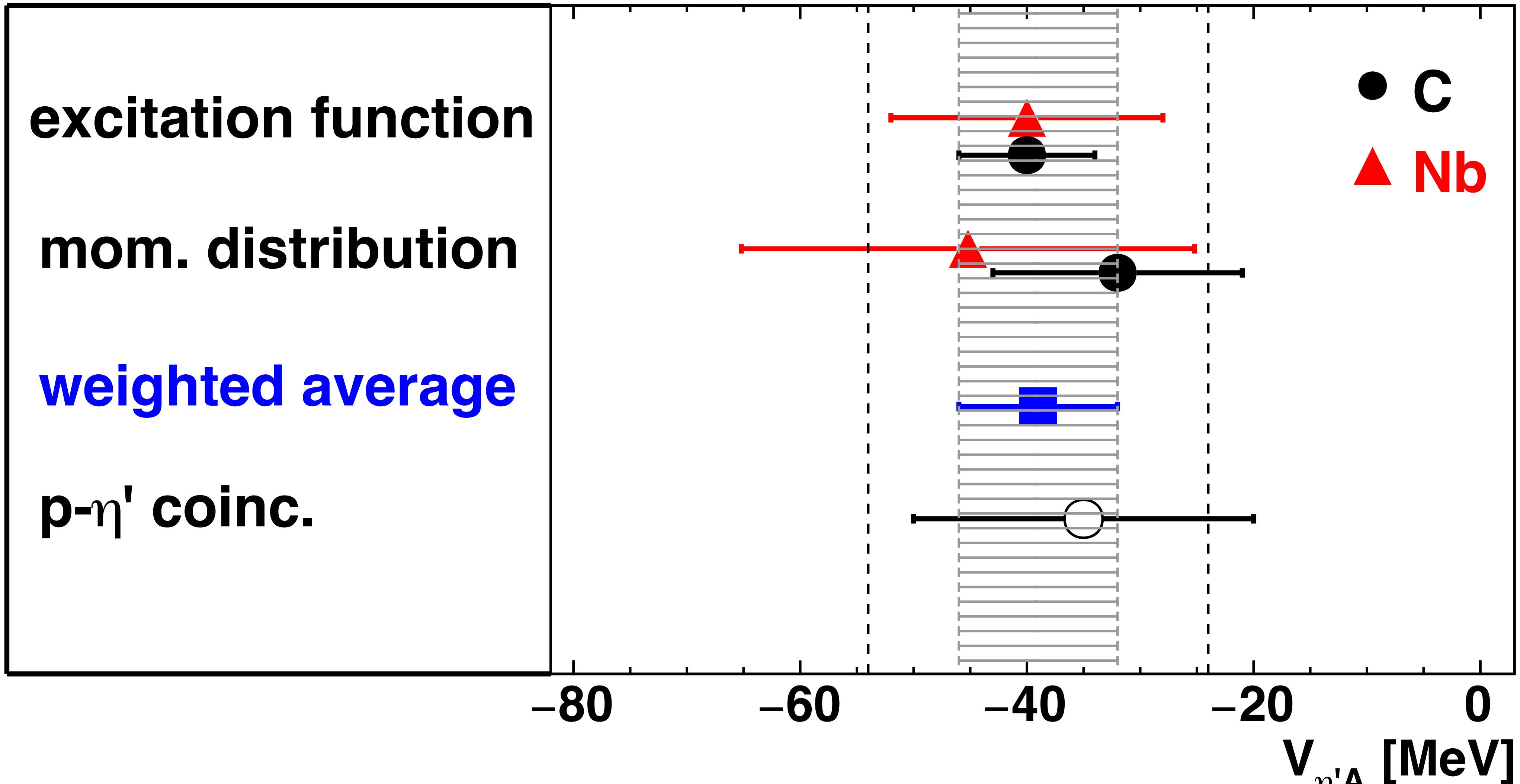
$V_{\eta'}(\rho=\rho_0) = -(40\pm6)$ MeV

$V_{\eta'}(<p_{\eta'}>\approx 1.1$ GeV/c$;\rho=\rho_0) = -(32\pm11)$ MeV

data disfavour strong mass shifts

determining the real part of the η'-nucleus potential

M. Nanova et al., PRC 94 (2016) 025205



$p_{\eta'} \lesssim m$

$p_{\eta'} \approx 600$ MeV/c

$V_0 = \Delta m(\rho=\rho_0) = -[39 \pm 7(stat) \pm 15(syst)]$ MeV

observed mass shift in agreement with QMC model predictions

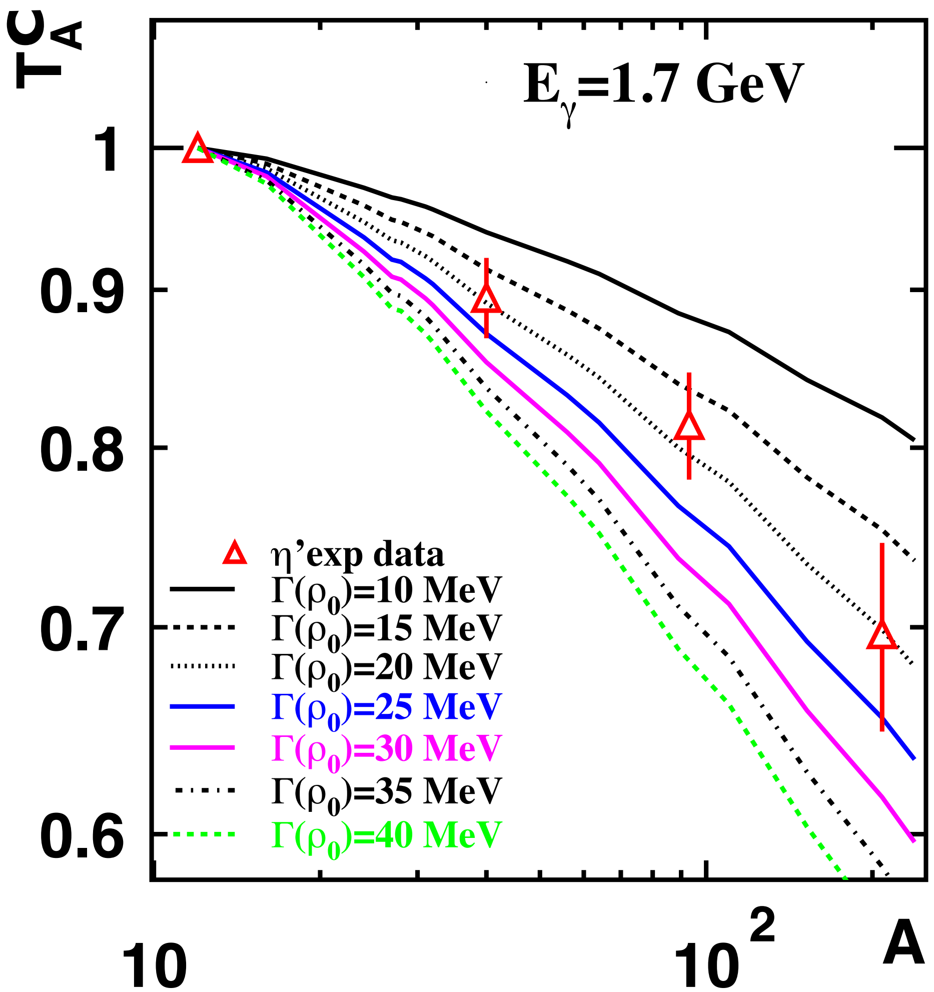
S. Bass and T. Thomas, PLB 634 (2006) 368

determining the imaginary part of the η'-nucleus potential

$$T_A = \frac{\sigma_{\gamma A \to \eta' X}}{A \cdot \sigma_{\gamma N \to \eta' X}}$$

mass dependence of T_A

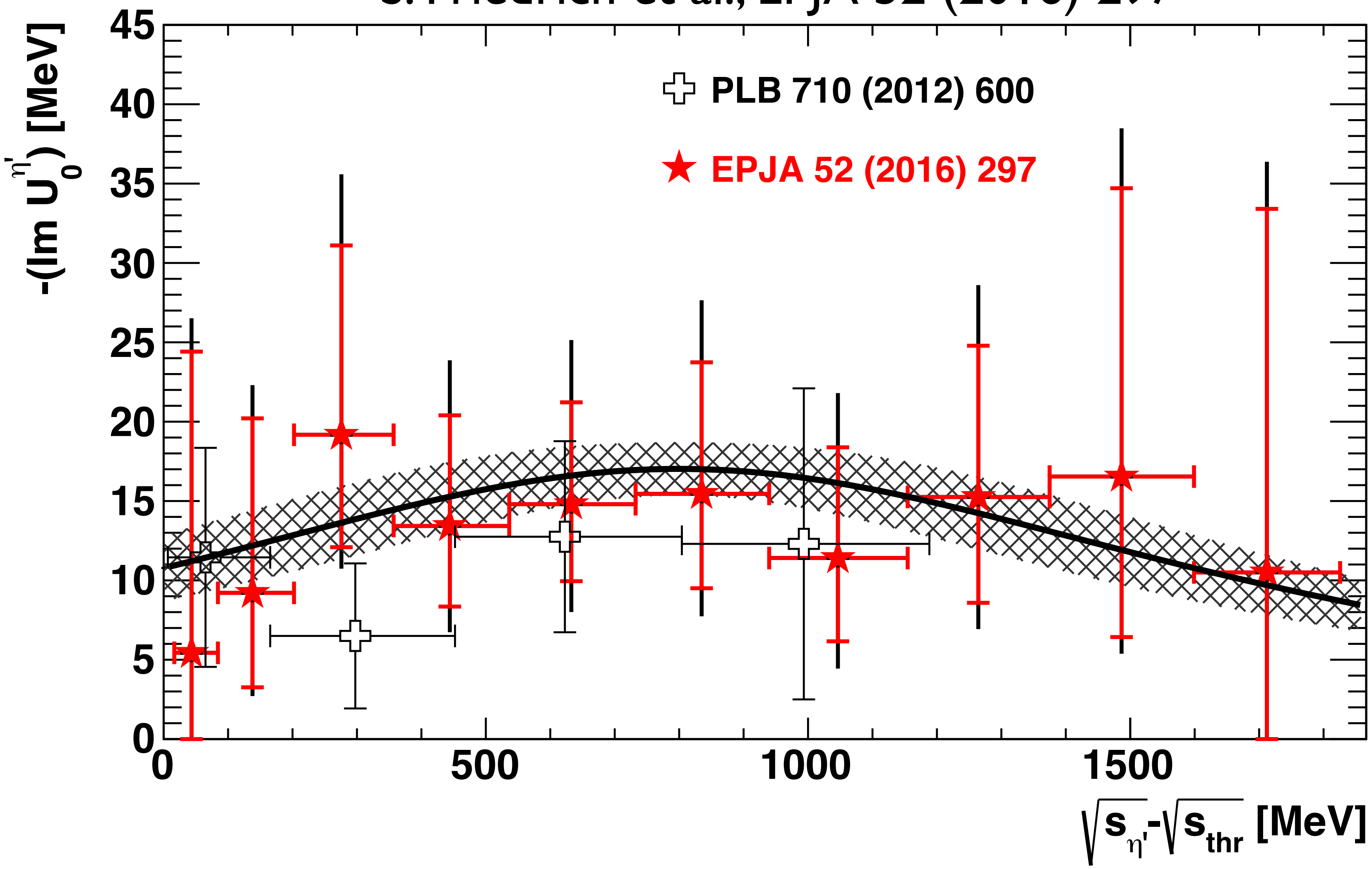
M. Nanova et al., PLB 710 (2012) 600



$\Gamma_{\eta'}(\rho=\rho_0)$ =15-25 MeV

momentum dependence of Γ_0, W_0

S. Friedrich et al., EPJA 52 (2016) 297



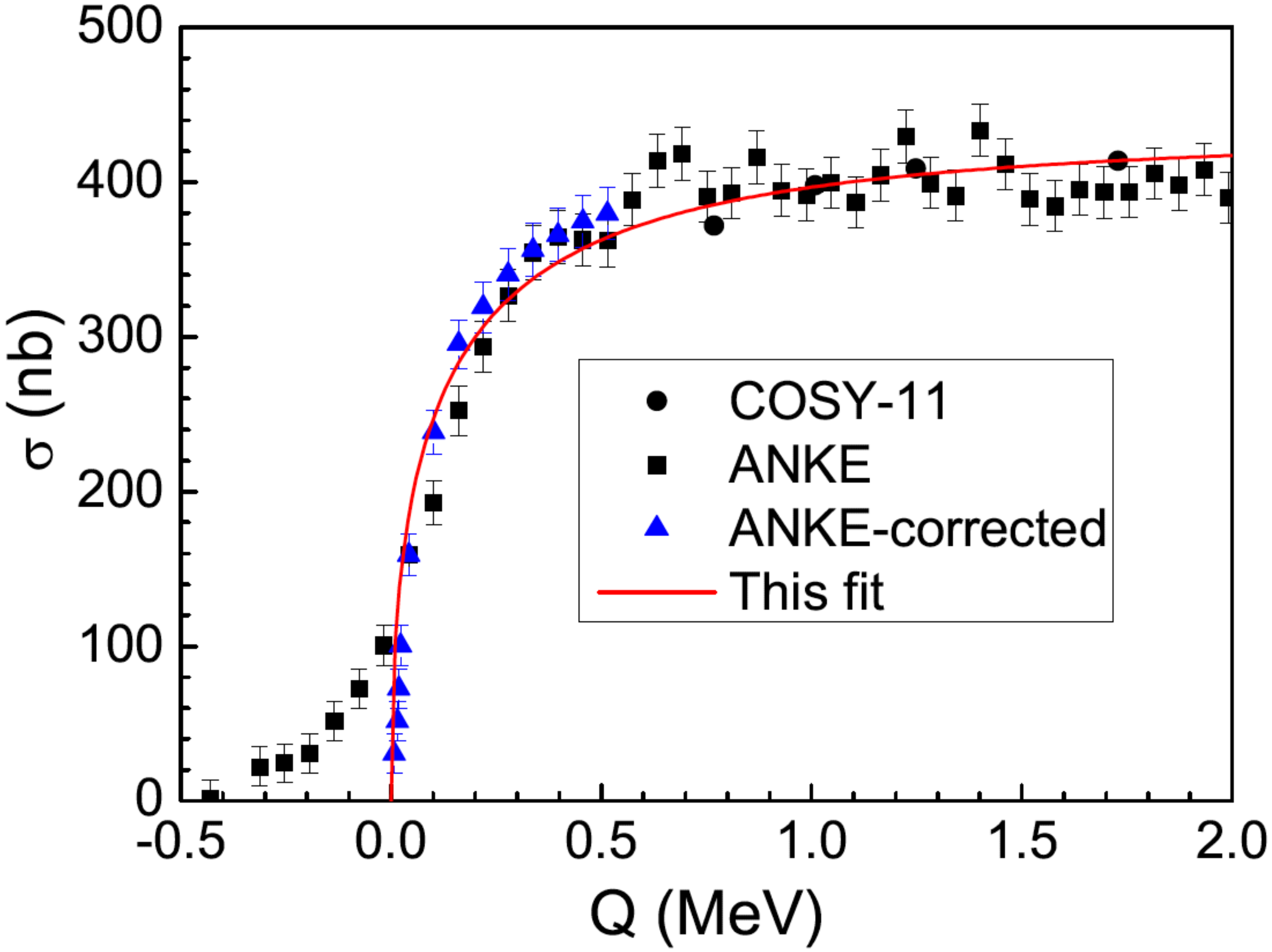
$$W_0 = \mathrm{Im}\,U(\rho=\rho_0, p_{\eta'}=0) = -[13 \pm 3(\mathrm{stat}) \pm 3(\mathrm{syst})]\ \mathrm{MeV}$$

real and imaginary part of the η-nucleus potential

$p+d \rightarrow \eta+{}^3He$

ANKE: T. Mersmann et al., PRL 98 (2007) 242301

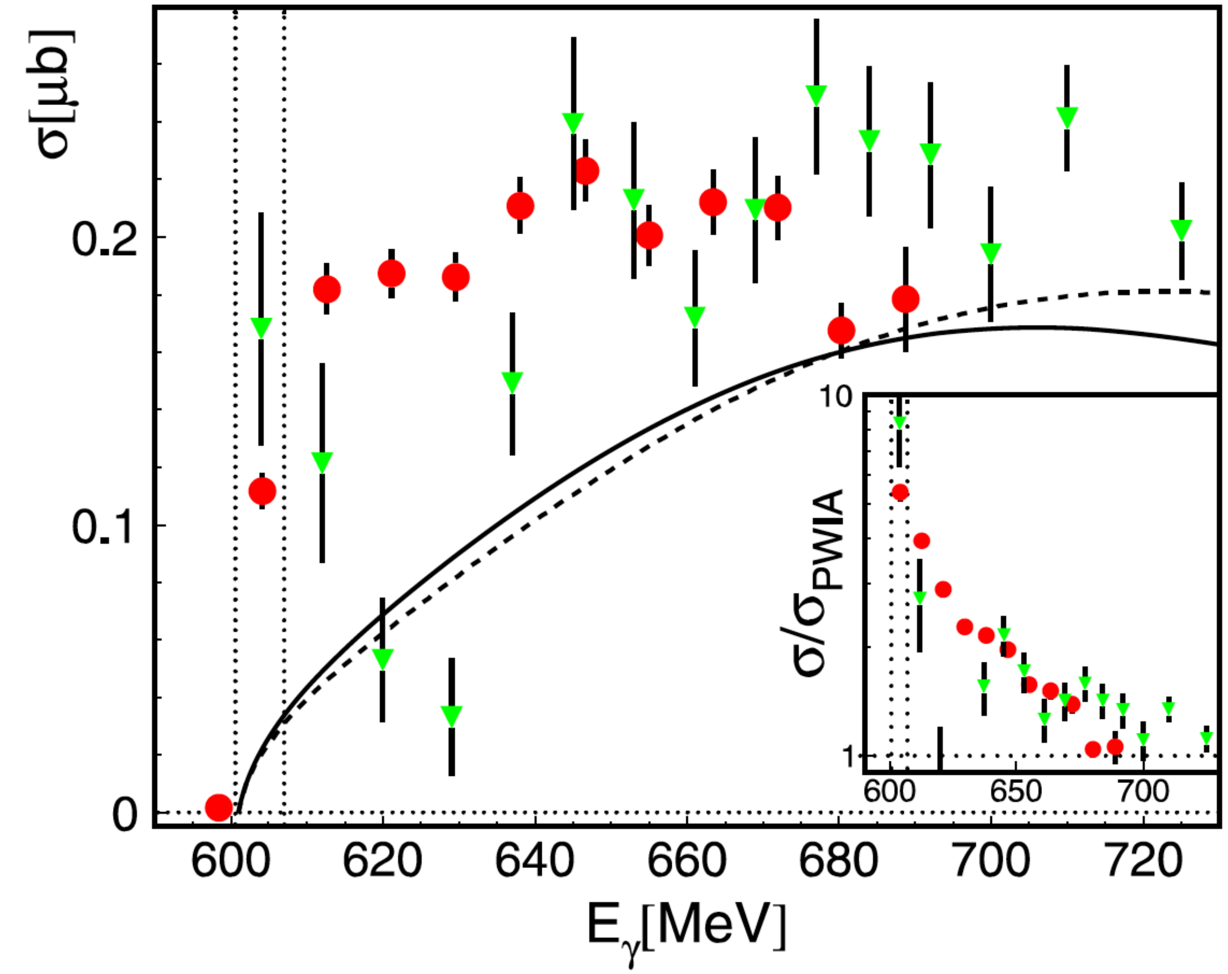
COSY-11: J. Smirski et al., PLB 649 (2007) 258



$\gamma+{}^3He \rightarrow \eta+{}^3He$

M. Pfeiffer et al., PRL 92 (2004) 252001

F. Pheron et al., PLB 709 (2012) 21



very steep rise of cross section near threshold !!

indication for a quasi-bound state near threshold ??

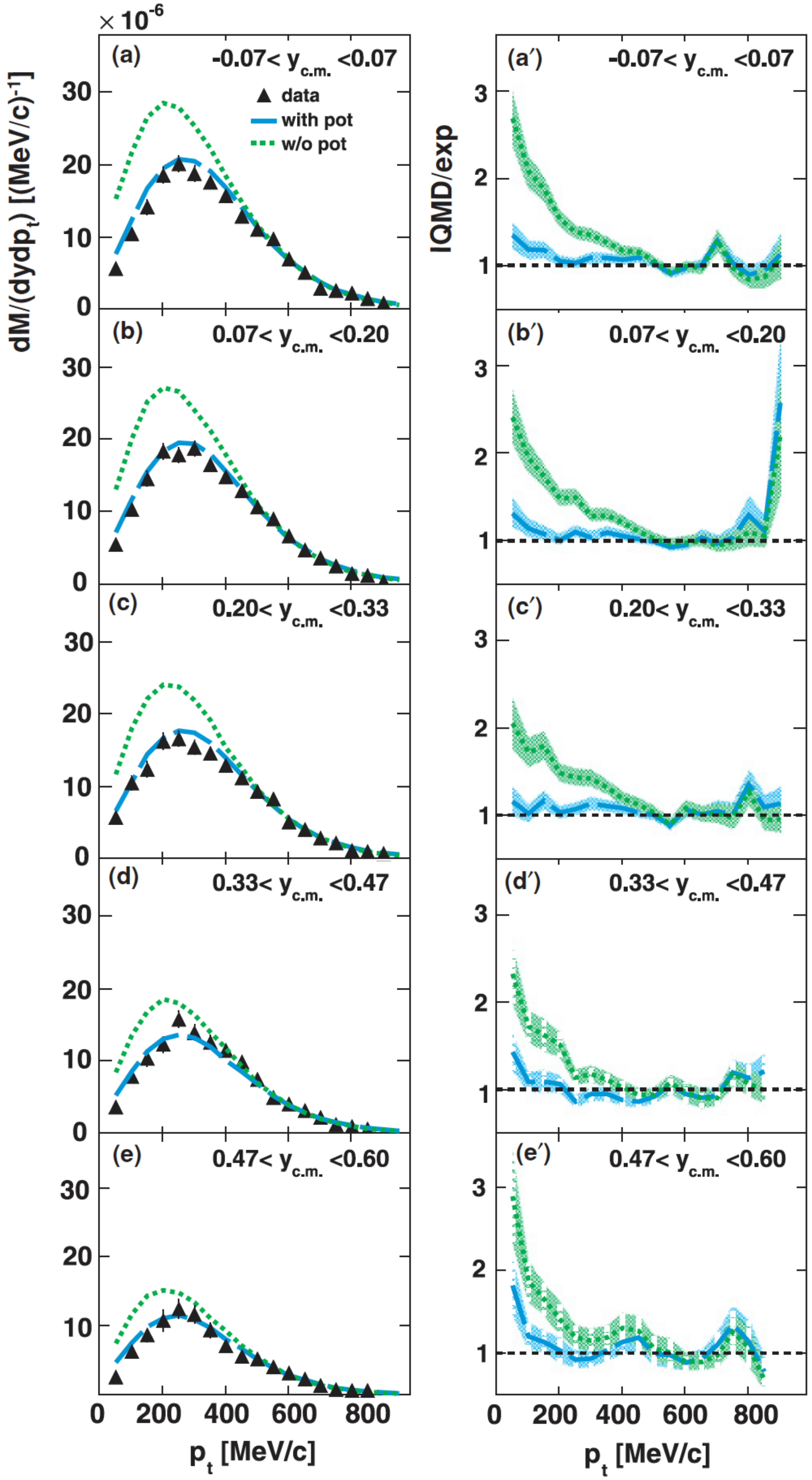
C. Wilkin et al. PLB 654 (2007) 92: pole at Q=-0.3 MeV; Γ = 0.3 MeV

J. J. Xie et al. PRC 95 (2017) 015202: BW structure at mass = -0.3 MeV; Γ = 3 MeV

V_0 = -(54±6) MeV; W_0 = -(20±2) MeV

⇒ talks by A. Gal, E. Oset and S. Hirenzaki, M. Skurzok (WASA)

determining the real part of the K^0 -nucleus potential



HADES: Ar + KCl at 1.756 AGeV

G. Agakishiev et al., PRC90 (2014) 054906

K^0 transverse momentum spectra
compared to IQMD transport calculations
without potential (dashed)
and with repulsive potential
of +46 MeV (solid curve)

$V \approx + 40$ MeV

⇒ talk by L. Fabbietti

information on the K$^-$-nucleus potential from atomic X-rays

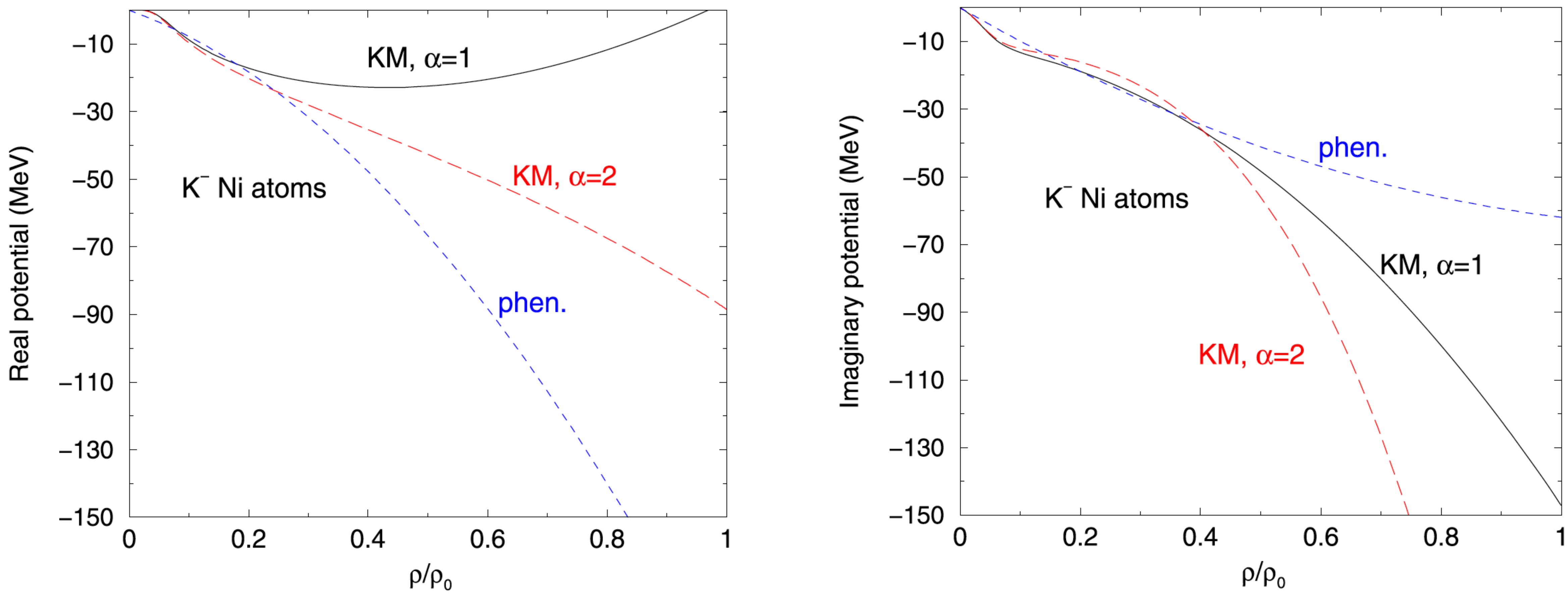
potential extracted from level shifts and widths of atomic states arising from the strong interaction

E. Friedman, A. Gal, Phys. Rep. 452 (2007) 89   S. Hirenzaki et al., PRC 61 (2000) 055205

X-rays from kaonic atoms sensitive to the potential at low nuclear densities

K$^-$ - multi nucleon absorption important (talks by C. Curceanu, J. Zmeskal)

E. Friedman, A. Gal, NPA 959 (2017) 66



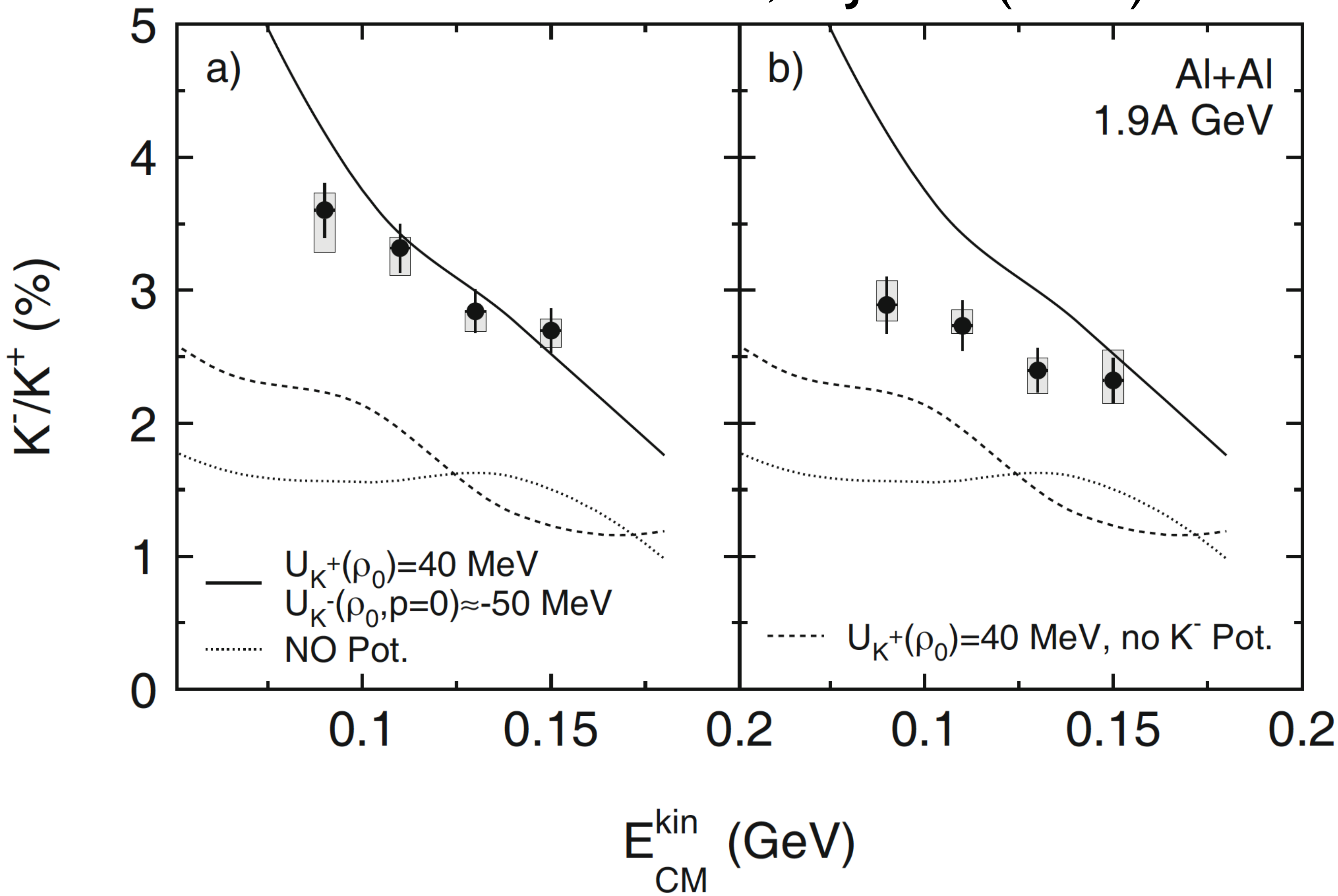
potentials can reliably be extracted up to 30% and 50% of normal nuclear matter density for real and imaginary part, respectively

other approaches needed to access $V(\rho=\rho_0), W(\rho=\rho_0)$

determining the real part of the K^--nucleus potential

K+ and K^- kinetic energy spectra from Al + Al at 1.94 AGeV

FOPI: P. Gasik et al., EPJA 52 (2016) 177



$V_{K^+} \approx +40$ MeV

$V_{K^-} \approx -50$ MeV ??

b.) corrected for feeding of K^- spectrum from decay $\Phi \rightarrow K^+K^-$ decays

Φ/K^--ratio = 0.36±0.05 Ni+Ni at 1.9 AGeV (FOPI)

Φ/K^--ratio = 0.52±0.16 Au+Au at 1.23 AGeV (HADES)

} not reproduced in transport calculations

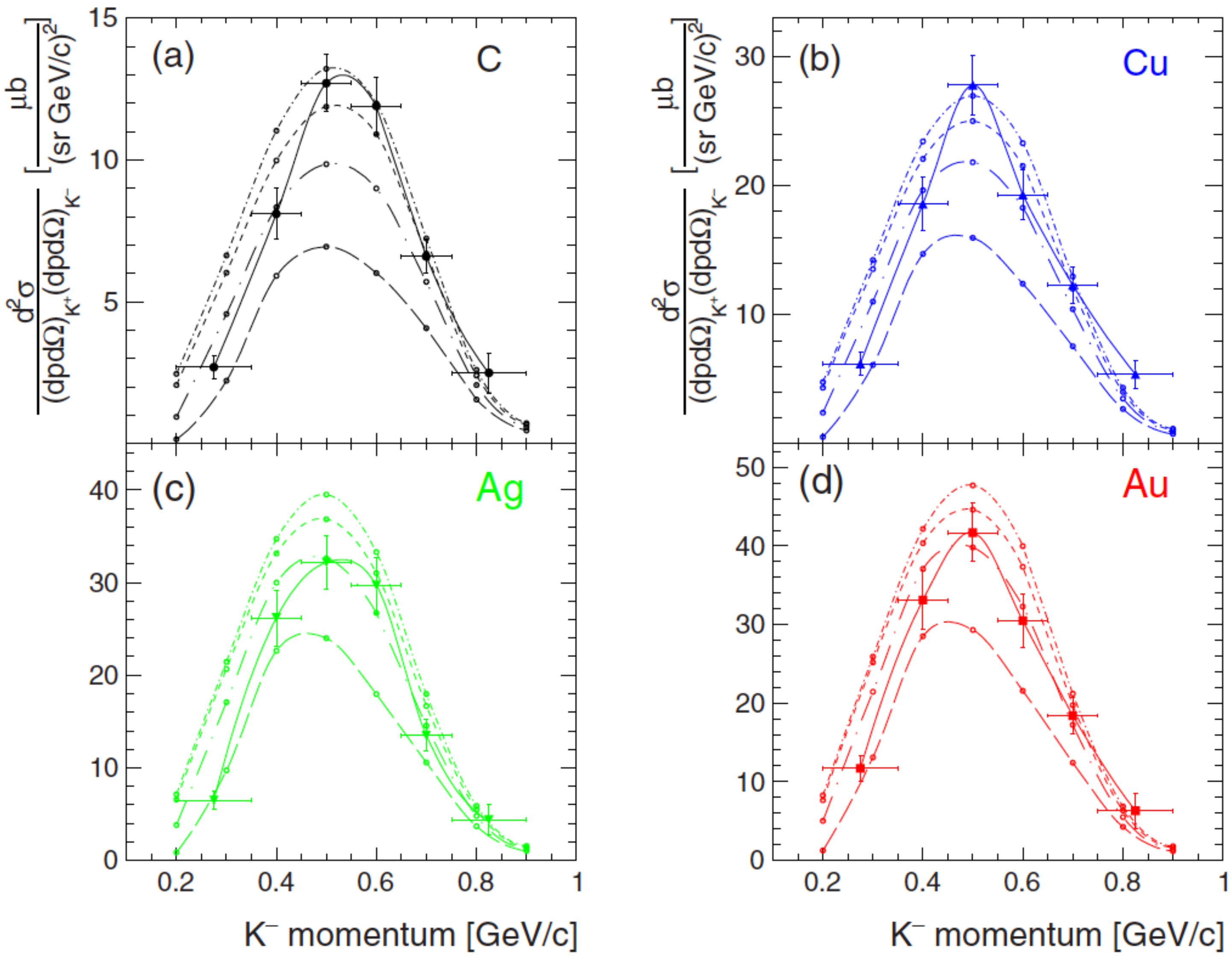
make sure other observables are reproduced before deducing potential parameters !!

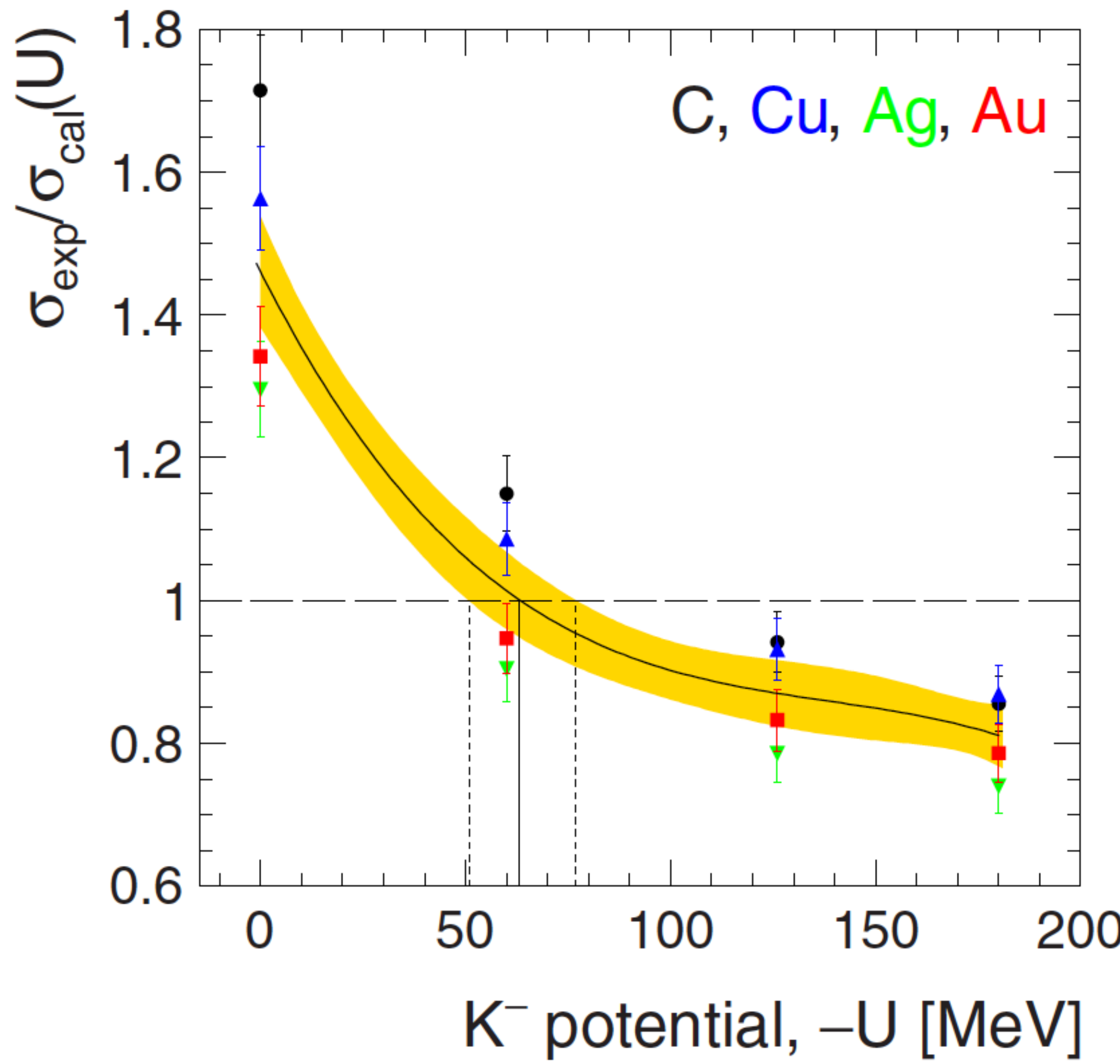
determining the real part of the K⁻-nucleus potential

$p + C, Cu, Ag, Au \rightarrow K^+ K^- + X$

$K^+ K^-$ - pairs not from Φ decay

ANKE: Yu. T. Kiselev et al., PRC92 (2015) 065201





K^--momentum spectra in coincidence with K^+ ($200 \leq p_{K+} \leq 600$ MeV/c) compared to collision model calculations: E. Paryev et al., J. Phys. G 42 (2015) 075107

$V_{K^-}(\rho=\rho_0) = -63^{+50}_{-30}$ MeV accounting for systematic uncertainties

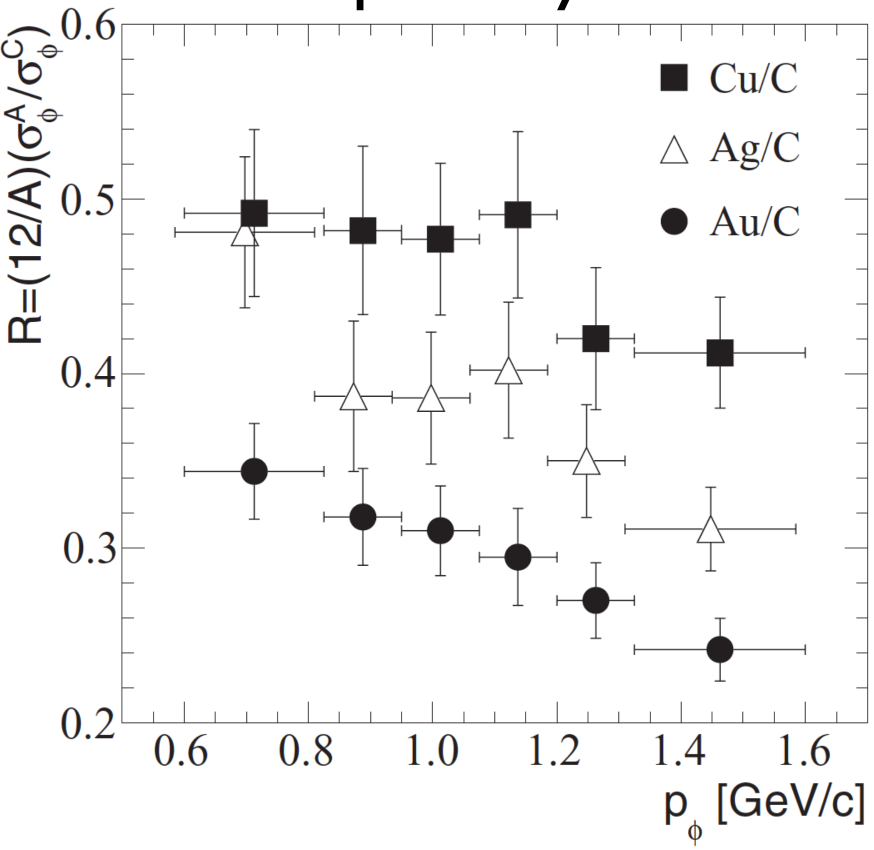
determining the imaginary part of the Φ-nucleus potential

M. Hartmann et al., PRC85 (2012)035206

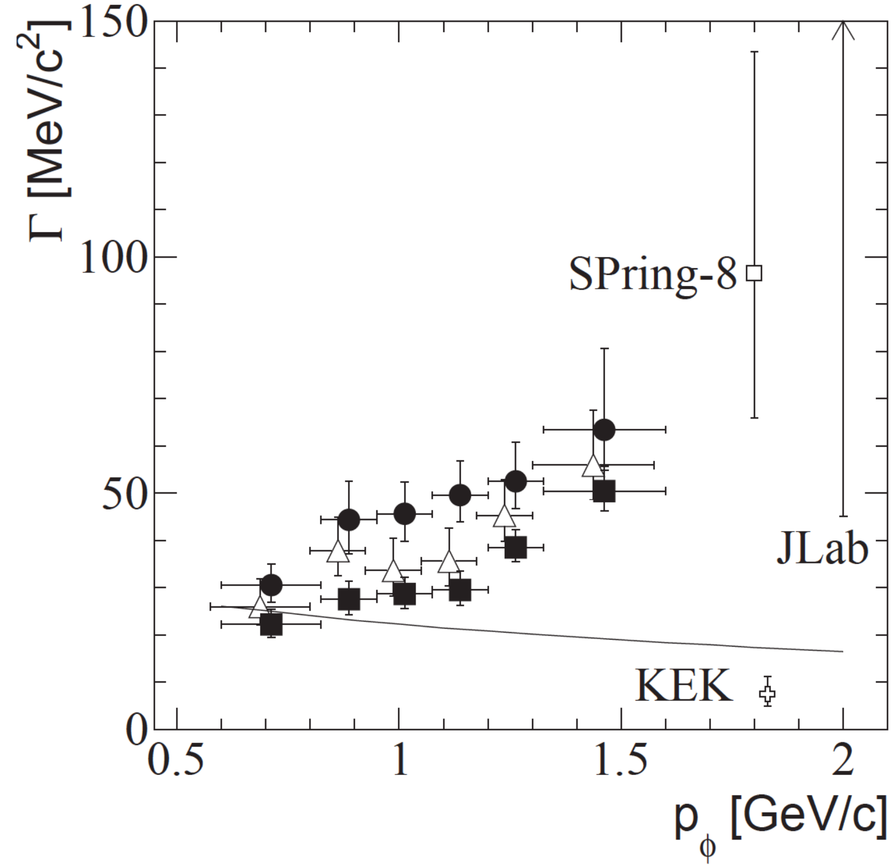
ANKE: p + C, Cu, Ag, Au → Φ + X at 2.83 GeV

momentum dependence of transparency ratio $T_A^C = \dfrac{\sigma_{\gamma A \to \phi X}}{A \cdot \sigma_{\gamma N \to \phi X}} \Big/ \dfrac{\sigma_{\gamma C \to \phi X}}{12 \cdot \sigma_{\gamma N \to \phi X}}$

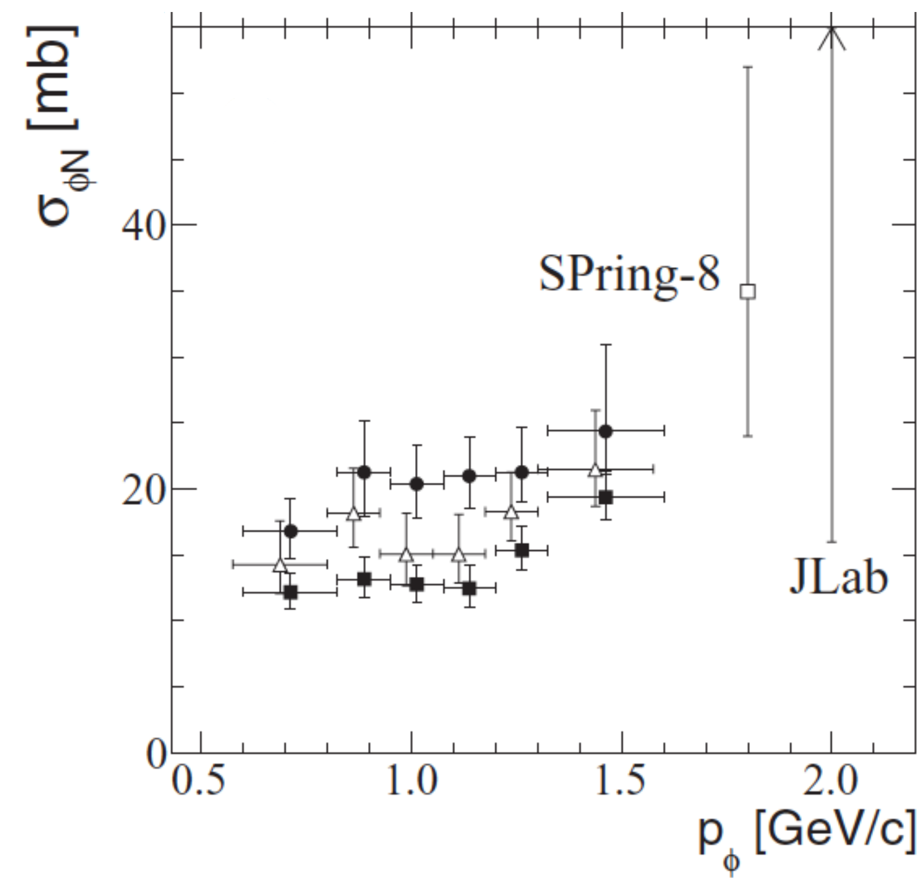
transparency ratio



in-medium width



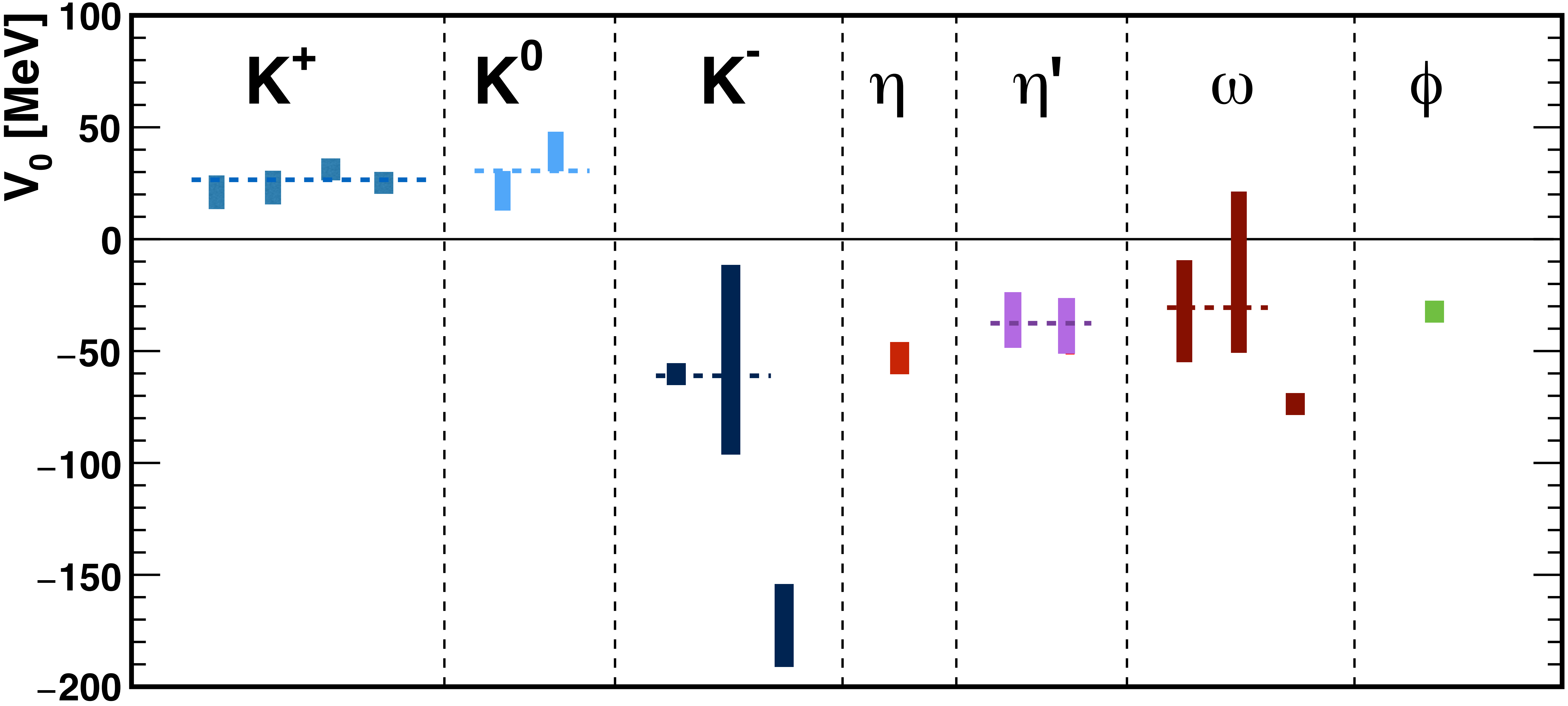
σ_{inel}



$$W(\rho=\rho_0) = -\Gamma/2\ (\rho=\rho_0) = -1/2 \cdot \hbar c \cdot \rho_0 \cdot \sigma_{inel} \cdot \beta$$

$$W = -\ (10\text{-}30)\ \text{MeV for } 0.7 < p_\Phi < 1.5\ \text{GeV/c}$$

real part of the meson-nucleus potential



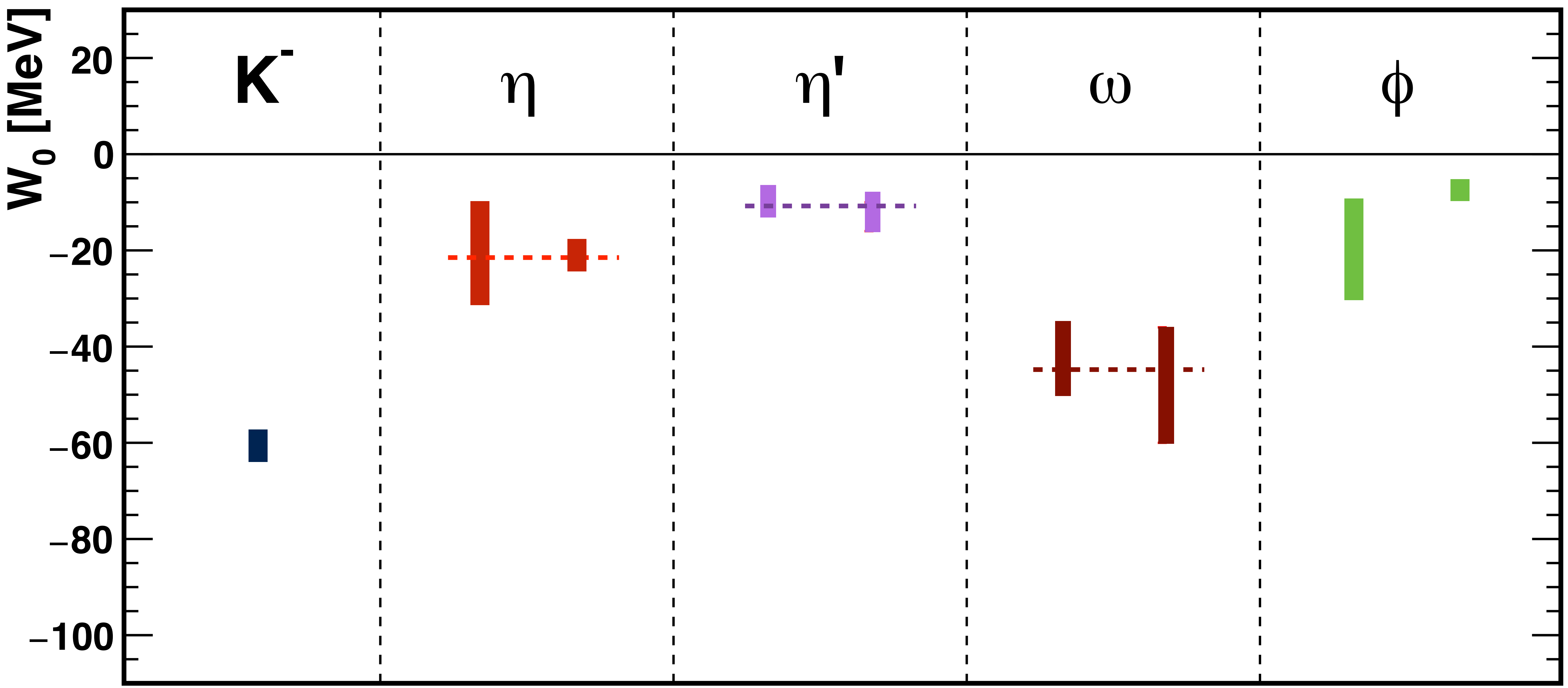
meson-nucleus real potential:

K^+, K^0 repulsive: 20-40 MeV

K^- strongest attraction: - (30 - 100) MeV

η, η', ω, Φ weakly attractive: - (20 - 50) MeV

imaginary part of the meson-nucleus potential



meson-nucleus imaginary potential:

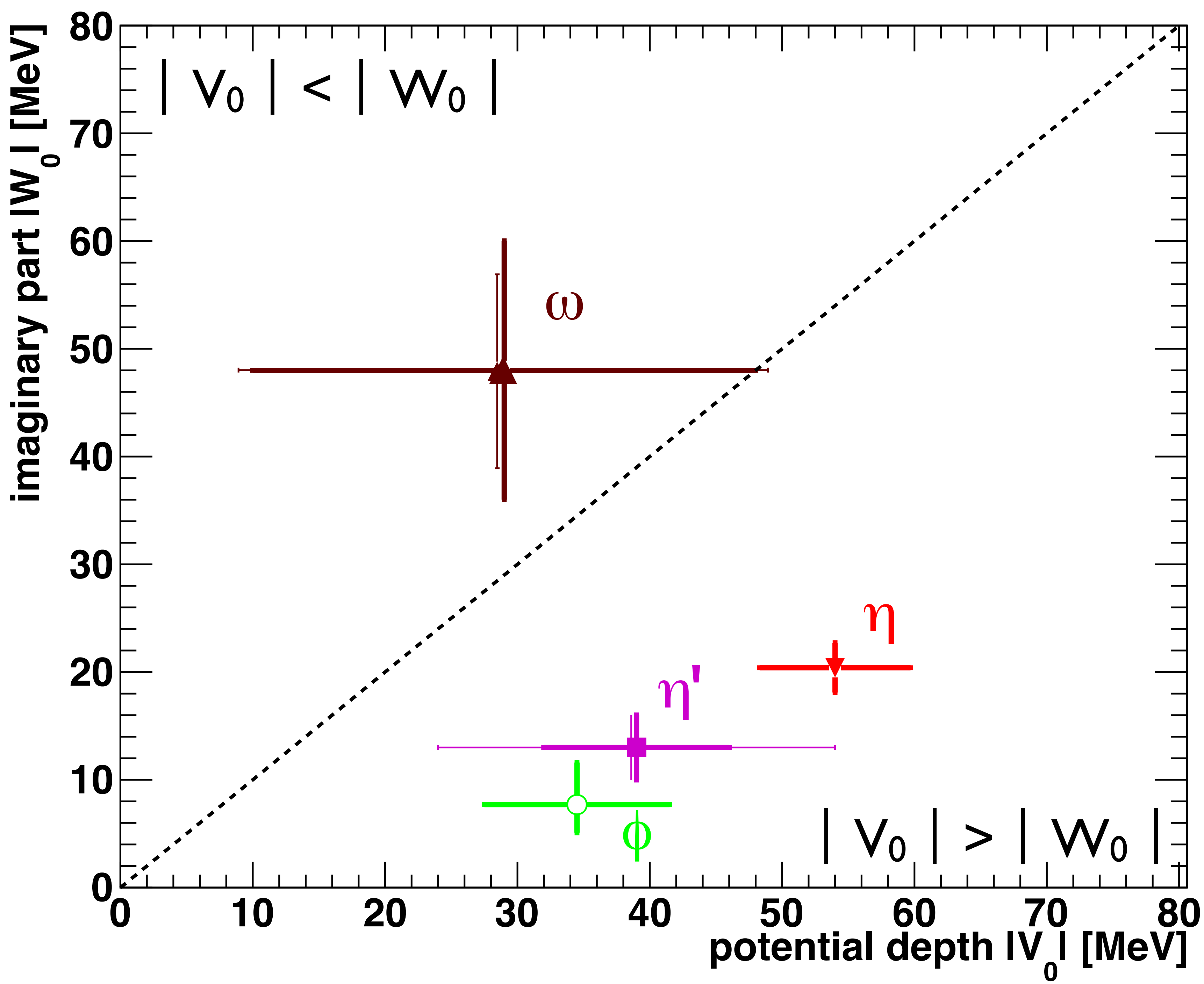
η' : ≈ -10 MeV

η, Φ : ≈ - 20 MeV

ω : ≈ - 40 MeV

K^- : ≈ - 60 MeV

real vs. imaginary part of the meson-nucleus potential



mesons with | V$_0$ | > | W$_0$ | suitable for search for meson-nucleus quasi-bound states

most favourable candidates: η, η'

search for meson-nucleus bound states

- K^- pp clusters (JPARC-E15): talk by T. Sekihara
- search for $K^- \otimes A$ bound states difficult because of large width due to K^--multi nucleon absorption — talks by K. Piscicchia, C.Curceanu
- search for $\eta \otimes {}^{3,4}He$ (WASA@COSY): talk by M. Skurzok
- search for $\eta' \otimes {}^{11}C$ (FRS@GSI): talk by Y. Tanaka

search for meson-nucleus bound states with Φ and heavier mesons (charm sector)

general experimental problem:

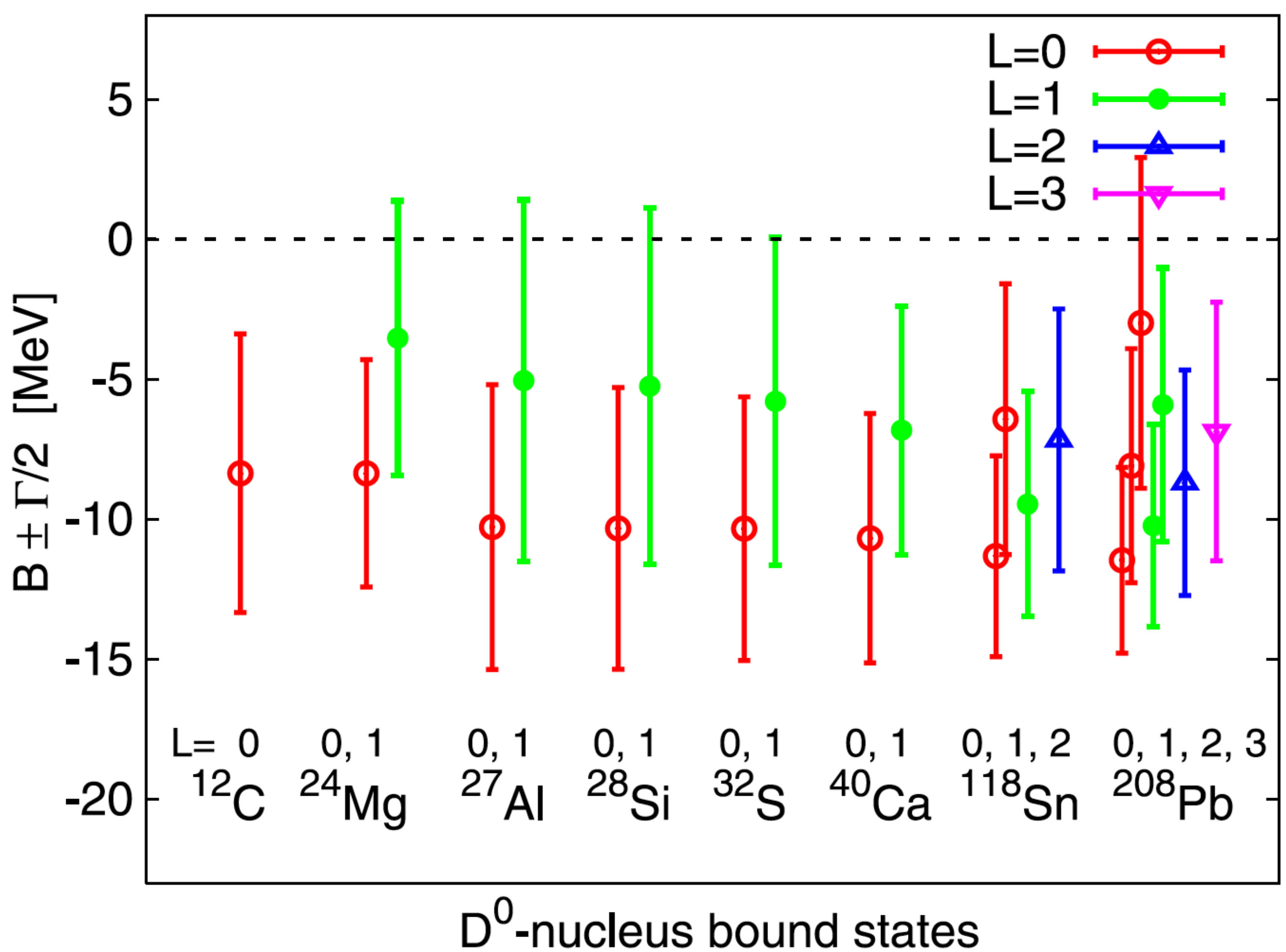
heavy meson production associated with high momentum transfer

probability for nucleus to stay intact ~ $F_A^2(q^2)$

minimising momentum transfer: M. Faessler, NPA 692 (2001) 104c

favourable reaction $\bar{p}\,p \rightarrow X\,Y$ with Y forward and X backward in cm

$$p_{min}(X) \approx \frac{m_X^2 - m_N^2}{2\,m_N}$$ (still 1.4 GeV/c for $\bar{D}D$ pairs !!)



more favourable:

two step production

$\bar{p}\,p \rightarrow D^{*-}\,D^{+}$

$D^{*-} + (Z,A) \rightarrow \pi^0 + D^{-} \otimes (Z,A)$

J. Yamagata-Sekihara et al., PLB 754 (1016) 26

summary and conclusions

- meson-nucleus interaction described by complex potential

$$U(r) = V(r) + i\,W(r)$$

- real part of meson-nucleus potential deduced from comparison of measured meson excitation functions or momentum distributions with transport and/or collision model calculations

- imaginary part of meson-nucleus potential deduced from comparison of measured transparency ratios with transport and/or collision model calculations

- measured potential parameters indicate favourable conditions ($|V_0| >> |W_0|$) for observing meson-nucleus quasi-bound states

  promising candidates: η, η'

- extension to charm sector difficult because of high momentum transfer