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# Investigation of the $^3\text{He}-\eta$ system with polarized beams at ANKE

## II International Symposium on Mesic Nuclei

September 22-25, 2013

wissen.leben  
WWU Münster

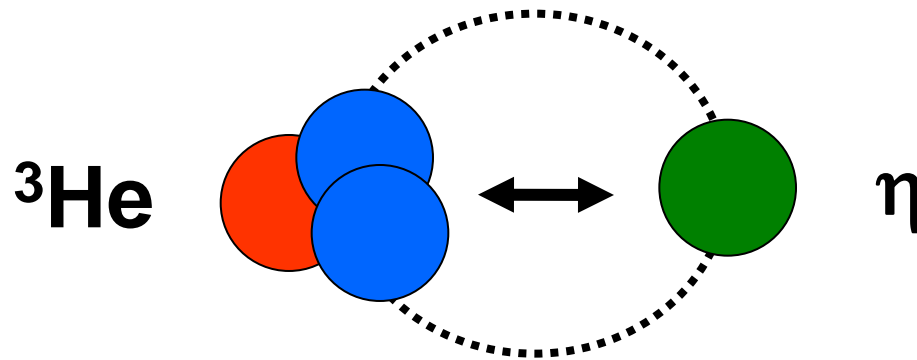
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## Why $\eta$ -Meson Production Close to Threshold?

- Do bound meson-nucleus systems exist?



- ANKE:  $d+p \xrightarrow{(\rightarrow)} ^3\text{He}+\eta$
- Excitation function close to threshold  $\rightarrow$  FSI
- Polarized beam  $\rightarrow$  Test of FSI hypothesis, role of spins

# The COSY-Accelerator at Jülich



## Energy range

- 0.045 – 2.8 GeV (p)
- 0.023 – 2.3 GeV (d)  
(momentum 3.7 GeV/c)

## Beam cooling

- Electron cooling
- Stochastic cooling

## Polarisation

- p, d beams & targets

## Beams

- internal, external

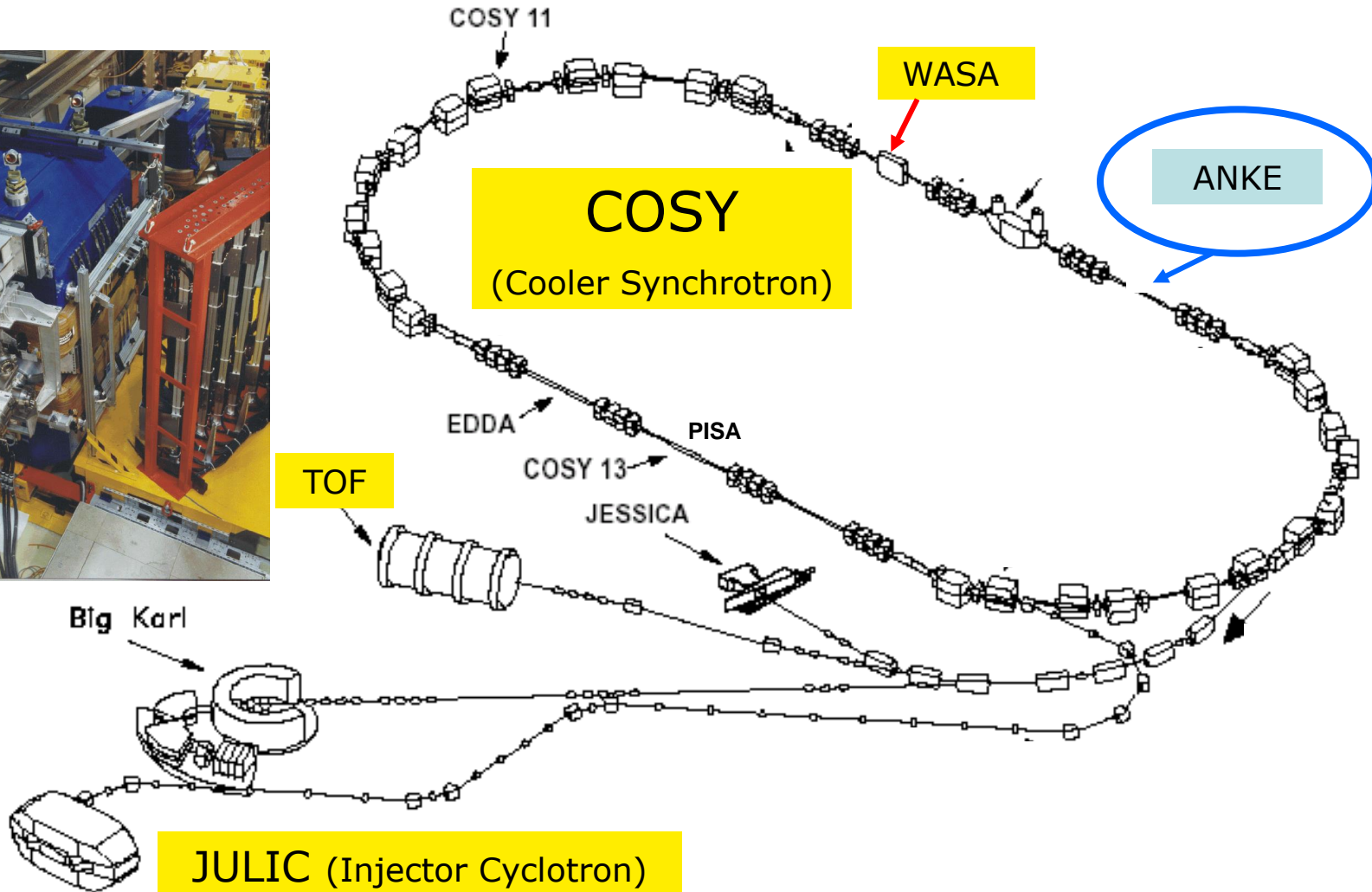
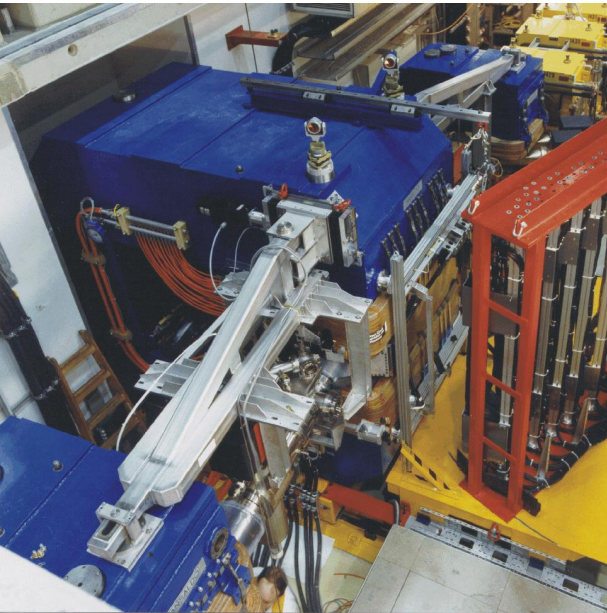
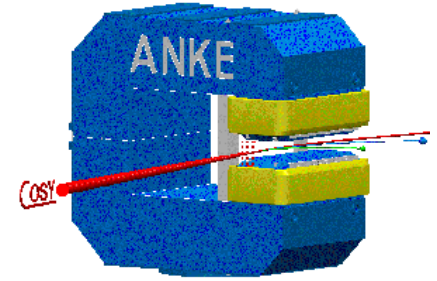
## Experiments, Detectors

- ANKE, TOF, WASA, ...

**COSY** (Cooler Synchrotron)

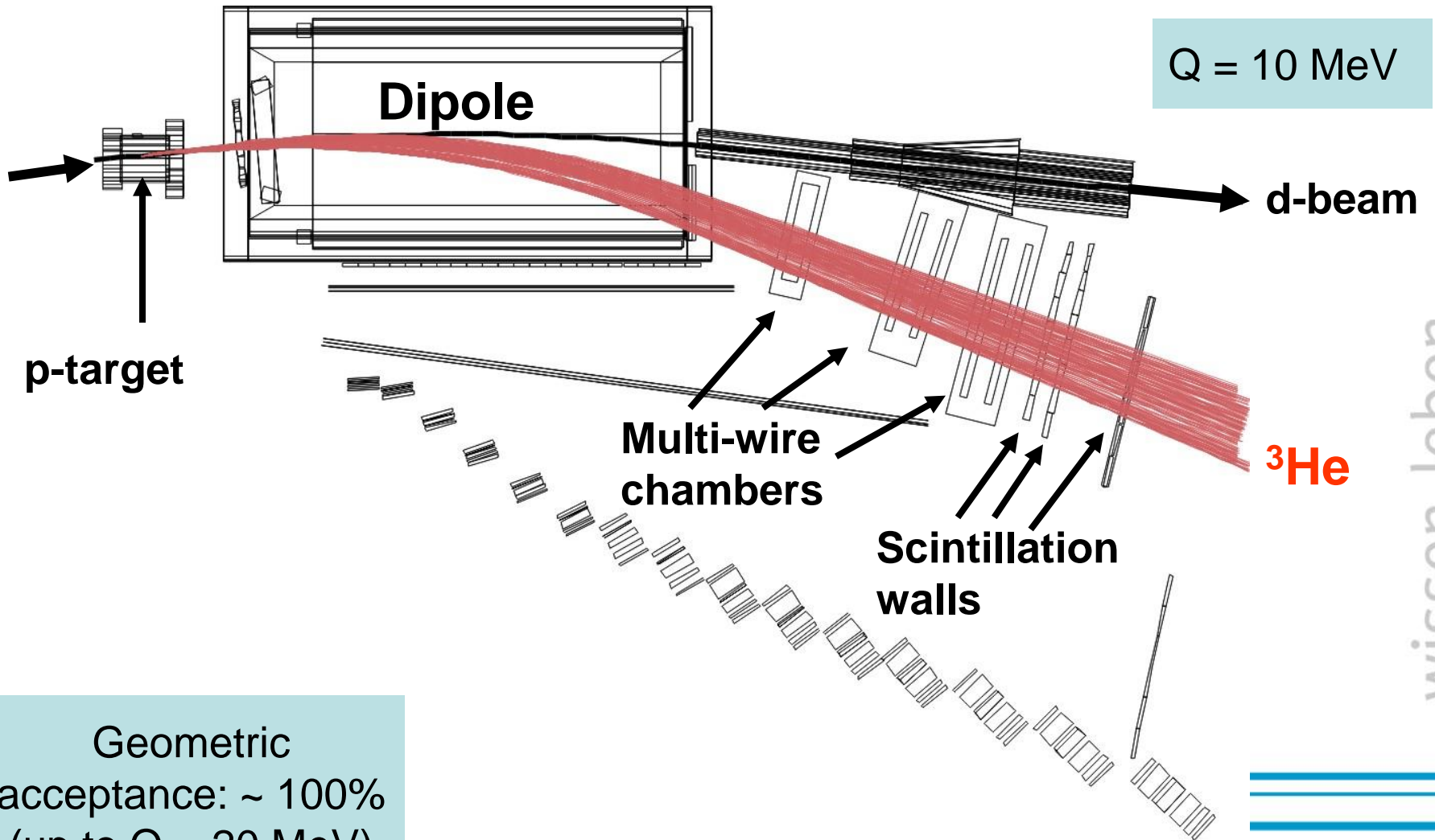


# The ANKE-Facility





# Identification of $^3\text{He}$ Nuclei at ANKE

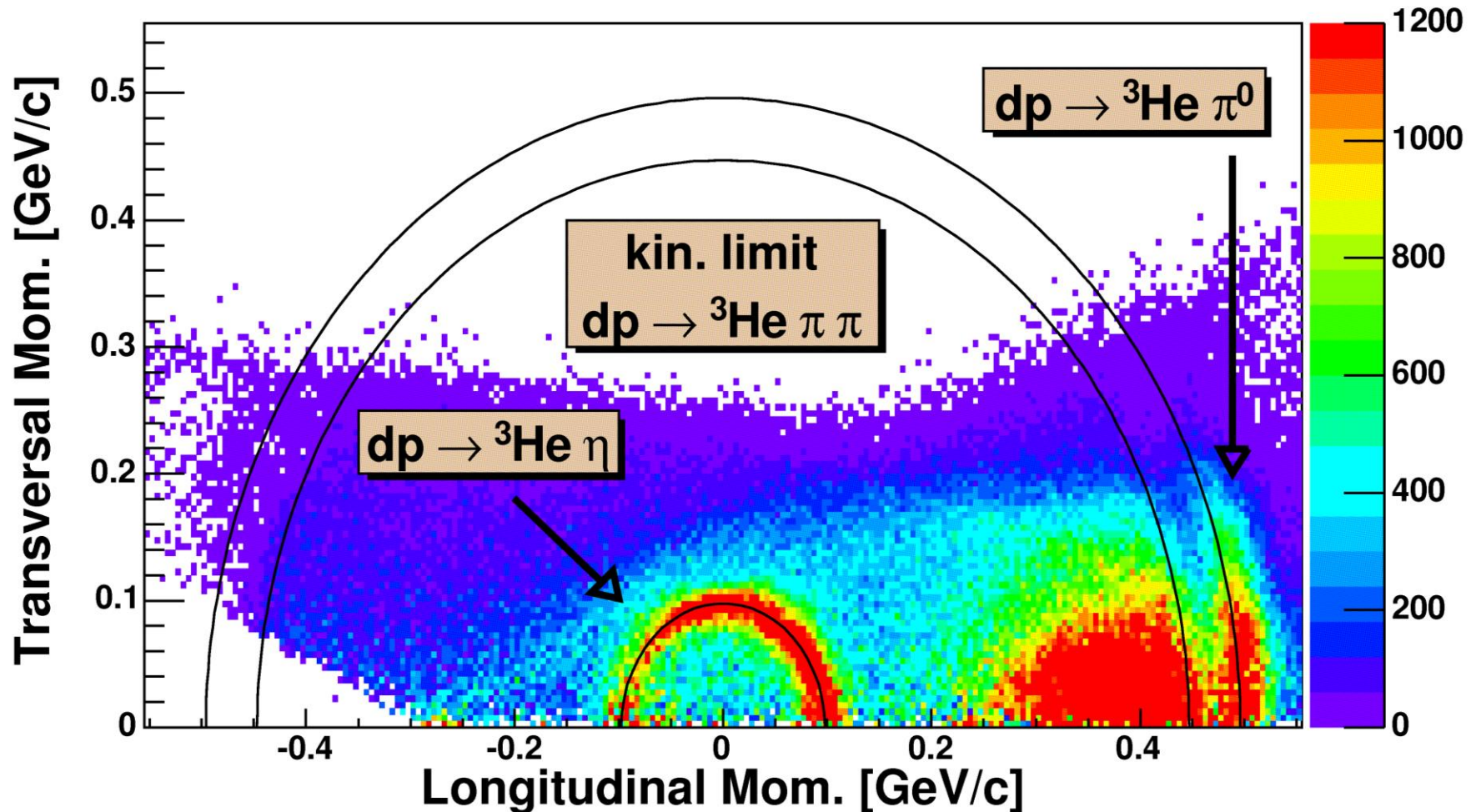


Geometric acceptance: ~ 100%  
(up to  $Q \sim 20$  MeV)



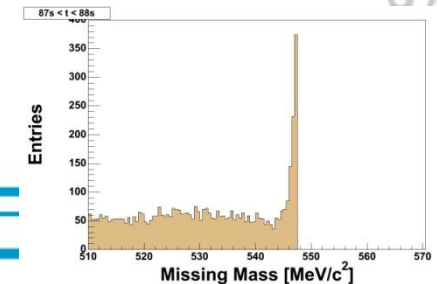
# Identification of the Reactions: $d+p \rightarrow {}^3\text{He}+X$

„Momentum rabbit“



## Identification of the Reactions: $d+p \rightarrow {}^3\text{He}+X$

- Energies and momenta of the incoming particles (d,p) known
  - Deuteron (mass =  $m_d$ ):  
energy + momentum: Adjustable by the accelerator
  - Proton (mass =  $m_p$ ):  
target particle at rest, momentum = 0
- Energy of the  ${}^3\text{He}$  nucleus measurable by detectors
- $\eta$ -meson: Not directly detectable at ANKE  
→ Identification of the reaction via the missing mass analysis



## Two-Particle Final State: Phase Space

Assumption:

- Two-particle reaction  $a+b \rightarrow c+d$  without initial and final state interactions („ISI“ and „FSI“):
- Scattering (and production) amplitude  $f = \text{const.}$ 
  - Increase of the cross section according to phase space expectations

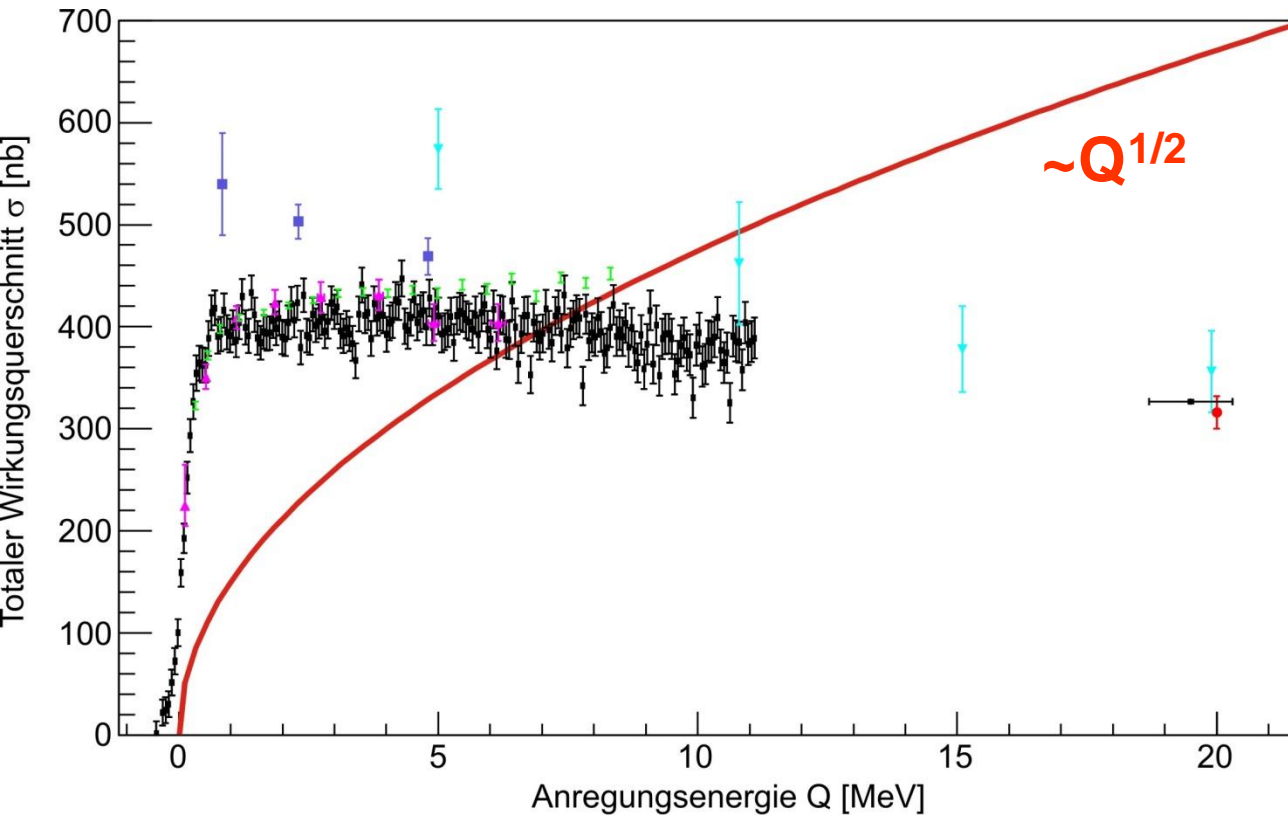
$$\frac{d\sigma(\mathcal{G})}{d\Omega} = \frac{p_f}{p_i} |f_s|^2 \propto p_f \propto \sqrt{Q}$$

$p_i / p_f$ : Momenta of in- and outgoing particles in the CMS

Q: Q-value = Sum of kinetic energies im CMS



## Results for the Reaction $d+p \rightarrow {}^3\text{He}+\eta$



But:

- Strong deviation from phase space expectation!
- Most probably not caused by higher partial waves



## The Reaction $d+p \rightarrow {}^3\text{He}+\eta$

- Extreme increase of the total cross section close to the production threshold
- Increase of the cross sections within  $\Delta Q < 1$  MeV
  - strong energy dependence at threshold
- After that total cross sections remain almost constant
  - Additional effect beside pure phase space

Explanation: Strong final state interaction (FSI) between  ${}^3\text{He}$  nucleus and  $\eta$ -meson

# Scattering Theory and Final State Interaction

Description of the cross section including FSI:

$$\frac{d\sigma(\mathcal{G})}{d\Omega} = \frac{p_f}{p_i} |f_s|^2 = \frac{p_f}{p_i} \cdot \frac{|f_{\text{prod}}|^2}{\left|1 - i \cdot a \cdot p_f + \frac{1}{2} a \cdot r_0 \cdot p_f^2\right|^2}$$

Assumption:

- Energy dependence of the production amplitude  $f_{\text{Prod}}$  is negligible close to threshold:  $f_{\text{Prod}} \sim \text{const.}$
- Initial State Interaction (ISI) also:  $\text{ISI} = \text{const.}$

# Scattering Theory and Final State Interaction

- The scattering length can deliver information about possible bound states
- Conditions for bound  $\eta^3\text{He}$  state:
  - Existence of a pole in the complex  $p_f$  plane

$$f_s = \frac{f_{\text{prod}}}{1 - i \cdot a \cdot p_f + \frac{1}{2} a \cdot r \cdot p_f^2}$$

$$a \equiv a_r + ia_i$$

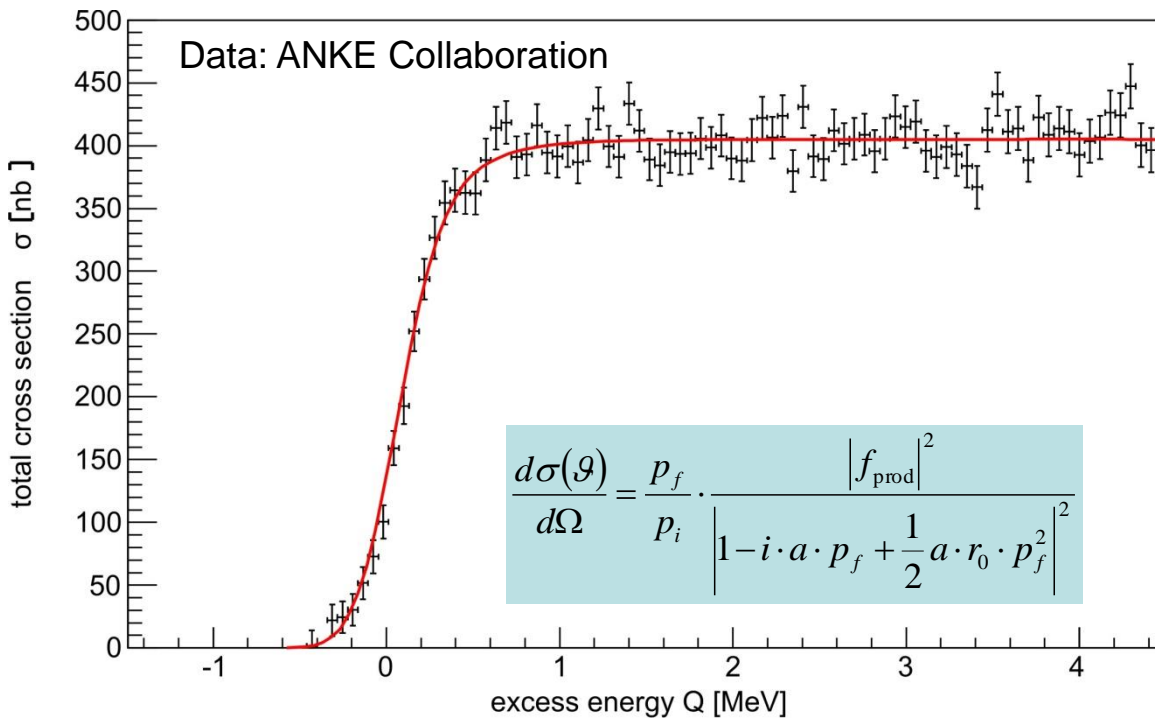
$$r \equiv r_r + ir_i$$

- As well as

$$a_r < 0, \quad a_i > 0, \quad R = \frac{|a_i|}{|a_r|} < 1$$

# The Reaction $d+p \rightarrow {}^3\text{He}+\eta$

Fit to data very close to threshold: Only s-wave



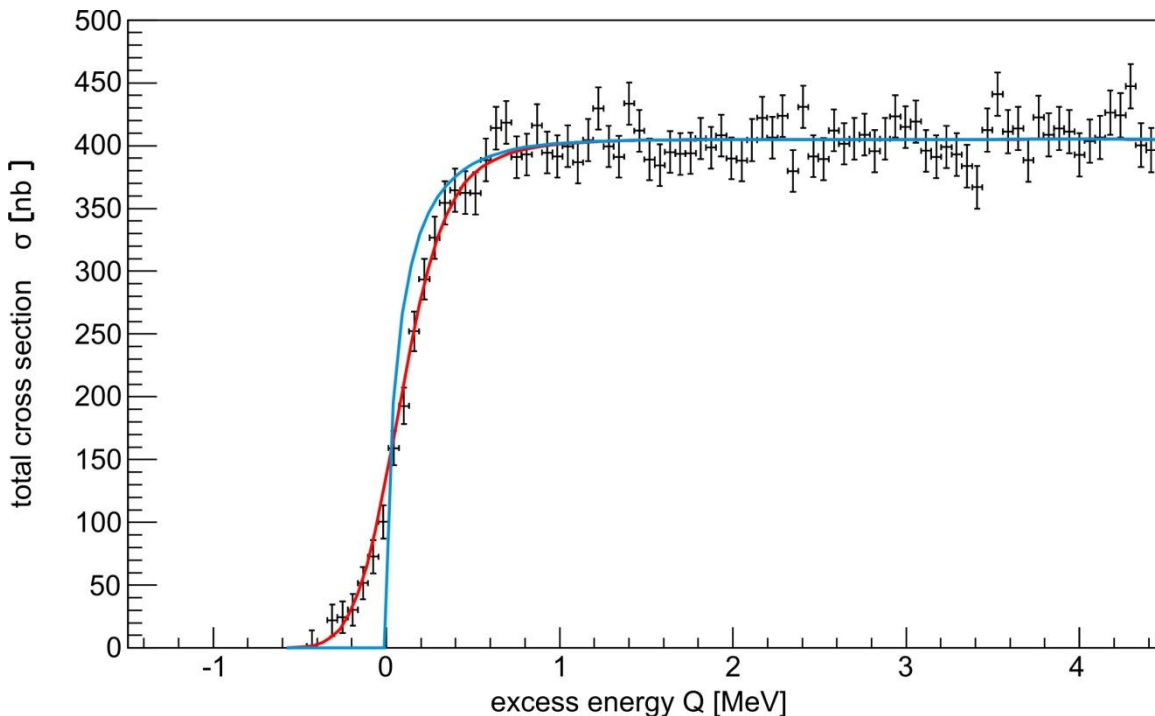
Fit parameter:

- Complex scattering length  $a = a_r + ia_i$
- Complex effective range  $r = r_r + ir_i$
- Finite momentum width  $\delta p_{\text{beam}}$  of the accelerator beam



# The Reaction $d+p \rightarrow {}^3\text{He}+\eta$

Excitation function without accelerator beam smearing  $\delta p_{\text{beam}}$ :



Blue line:

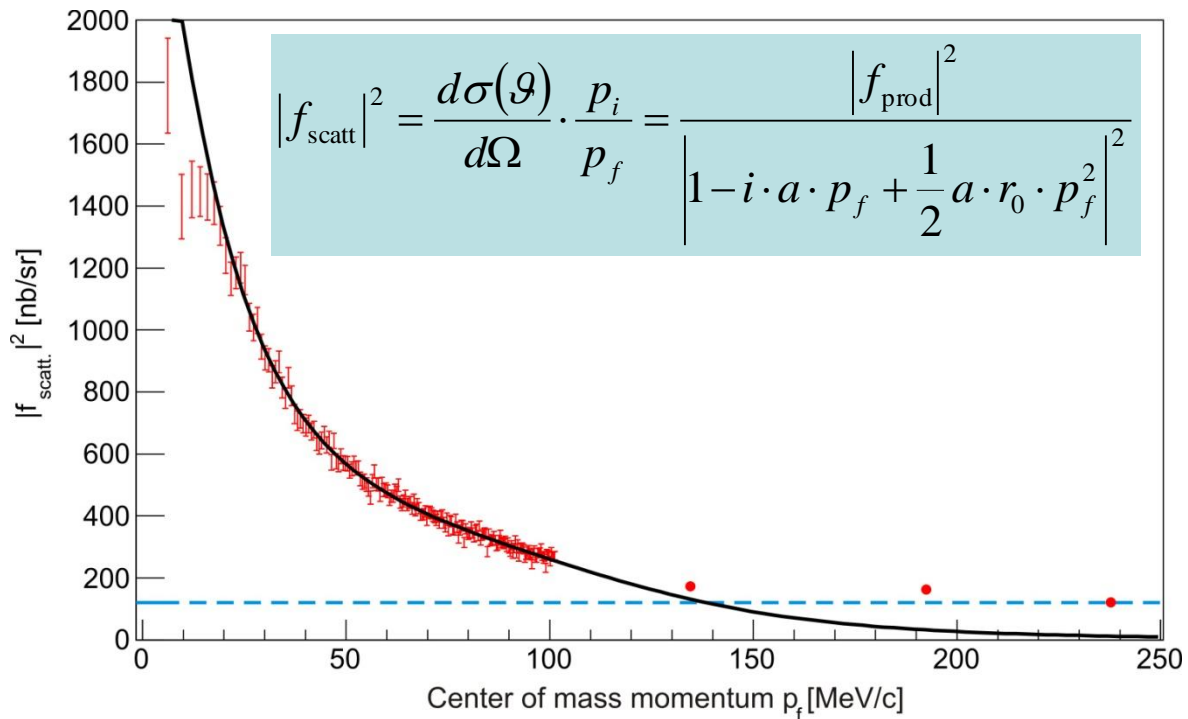
- Defolded shape, extracted from data (no accelerator beam smearing)

→

- Total cross section reaches maximum already  $\Delta Q < 0.5$  MeV above threshold

# The $d+p \rightarrow {}^3\text{He}+\eta$ Scattering Amplitude

Extracted scattering amplitude ( $Q > 0$  MeV)



- Scattering amplitude decreases rapidly with increasing final state momentum  $p_f$
- Scattering amplitude almost constant at high energies

→ strong FSI in  $\eta^3\text{He}$  system

## $\eta$ - $^3\text{He}$ Scattering Length

Fit to data delivers information about the complex  $\eta$ - $^3\text{He}$  scattering length:

$$\left( \frac{d\sigma(\mathcal{G})}{d\Omega} \right) \cdot \frac{p_i}{p_f} = |f_{\text{scat}}|^2 = |f_{\text{prod}} \cdot FSI|^2 = |f_{\text{prod}}|^2 \cdot |FSI|^2$$

Result:

$$a = \left[ \pm \left( 10.7 \pm 0.8_{-0.5}^{+0.1} \right) + i \left( 1.5 \pm 2.6_{-0.9}^{+1.0} \right) \right] \text{fm}$$

$$FSI = \frac{1}{1 - i \cdot a \cdot p_f + \frac{1}{2} a \cdot r_0 \cdot p_f^2}$$

Notice: Determination of  $|a_r|$ !



## $\eta$ - $^3\text{He}$ -Interaction: Determination of Poles

$$\left( \frac{d\sigma(\mathcal{G})}{d\Omega} \right) \cdot \frac{p_i}{p_f} = |f_{\text{scatt}}|^2 = |f_{\text{prod}} \cdot FSI|^2 = |f_{\text{prod}}|^2 \cdot |FSI|^2$$

$$FSI = \frac{1}{1 - i \cdot a \cdot p_f + \frac{1}{2} a \cdot r_0 \cdot p_f^2}$$

$$FSI = \frac{1}{\left(1 - \frac{p_f}{p_1}\right) \cdot \left(1 - \frac{p_f}{p_2}\right)}$$

$$a = -i \cdot \frac{p_1 + p_2}{p_1 \cdot p_2} \quad r_0 = + \frac{2 \cdot i}{p_1 + p_2}$$

$$p_1 = \left[ (-5 \pm 7_{-1}^{+2}) \pm i \cdot (19 \pm 2 \pm 1) \right] \text{MeV/c}$$

$$p_2 = \left[ (106 \pm 5) \pm i \cdot (76 \pm 13_{-2}^{+1}) \right] \text{MeV/c}$$

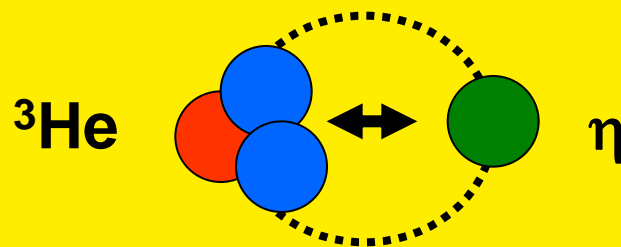
## $\eta$ - $^3\text{He}$ -Interaction: Determination of Poles

- Pole close to the reaction threshold

$$|Q_0| = \left| \frac{p_1^2}{2 \cdot m_{red}} \right| = 0.37 \text{ MeV}$$

- Position of the near-threshold pole (and scattering length) stable, i.e. nearly independent of fit range
- Large real part of scattering length and  $|a_r| > a_i$

→ indication for the existence of a bound state



(strong interaction!)



# Polarized Measurements

Production amplitude for  $dp \rightarrow {}^3\text{He} + \eta$  ( $\pi^0$ ):

$$f_B = \bar{u}_\tau \vec{p}_p \cdot (A\vec{\varepsilon}_d + iB\vec{\varepsilon}_d \times \vec{\sigma}) u_p$$

see:  
C. Kerboul et al.,  
Phys. Lett. B 181, 28 (1986)

Determination of the  
energy dependence  
of the amplitudes **A**  
and **B** by measurement  
of:

$$\frac{d\sigma}{d\Omega} = \frac{1}{3} \frac{p_\eta}{p_p} \left[ |A|^2 + 2|B|^2 \right]$$

$$T_{20} = \sqrt{2} \left[ \frac{|B|^2 - |A|^2}{|A|^2 + 2|B|^2} \right]$$

$$|A|^2 = \frac{p_p}{p_\eta} (1 - \sqrt{2}T_{20}) \frac{d\sigma}{d\Omega}$$

$$|B|^2 = \frac{p_p}{p_\eta} \left( 1 + \frac{1}{\sqrt{2}} T_{20} \right) \frac{d\sigma}{d\Omega}$$

$$T_{20} = \frac{2 \cdot \sqrt{2}}{p_{zz}} \cdot \frac{d\sigma_0 / d\Omega(\mathcal{G}) - d\sigma_\uparrow / d\Omega(\mathcal{G})}{d\sigma_0 / d\Omega(\mathcal{G})}$$

$$\mathcal{G} = 0^\circ \text{ or } 180^\circ$$

## Polarized Measurements

Assumption:  $\vec{d}p \rightarrow {}^3\text{He} + \eta$

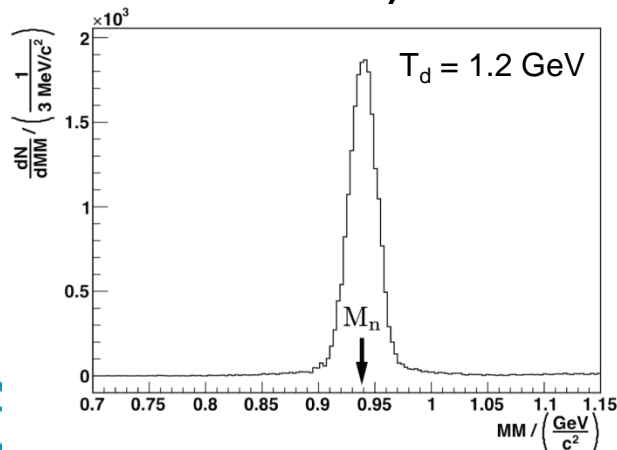
- Negligible effect of ISI
- Energy dependence of  $|f|^2$  only given by FSI
  - Shape of excitation function independent of spins
  - Same energy dependence of amplitudes  $|A|^2$  and  $|B|^2$

$$\begin{aligned} |A|^2 &= |A_0|^2 \cdot FSI(p_\eta) \\ |B|^2 &= |B_0|^2 \cdot FSI(p_\eta) \end{aligned} \quad \Rightarrow \quad T_{20} = \sqrt{2} \left[ \frac{|B_0|^2 - |A_0|^2}{|A_0|^2 + 2|B_0|^2} \right] \cdot \frac{FSI(p_\eta)}{FSI(p_\eta)} = \text{const.}$$

- Measure  $T_{20}$  as function of the excess energy

## The Reaction $d+p \rightarrow {}^3\text{He}+\eta$ at ANKE

- Alternating injection of unpolarized and tensor polarized deuterons in COSY
- Ramped COSY beam:  $Q = -5 \text{ MeV} \dots +10 \text{ MeV}$  (300 s)
- Full geometrical acceptance of ANKE for  $d+p \rightarrow {}^3\text{He}+\eta$
- Determination of  $p_{zz}$  by, e.g.,  $d+p \rightarrow (pp)+n$  (analyzing powers known)

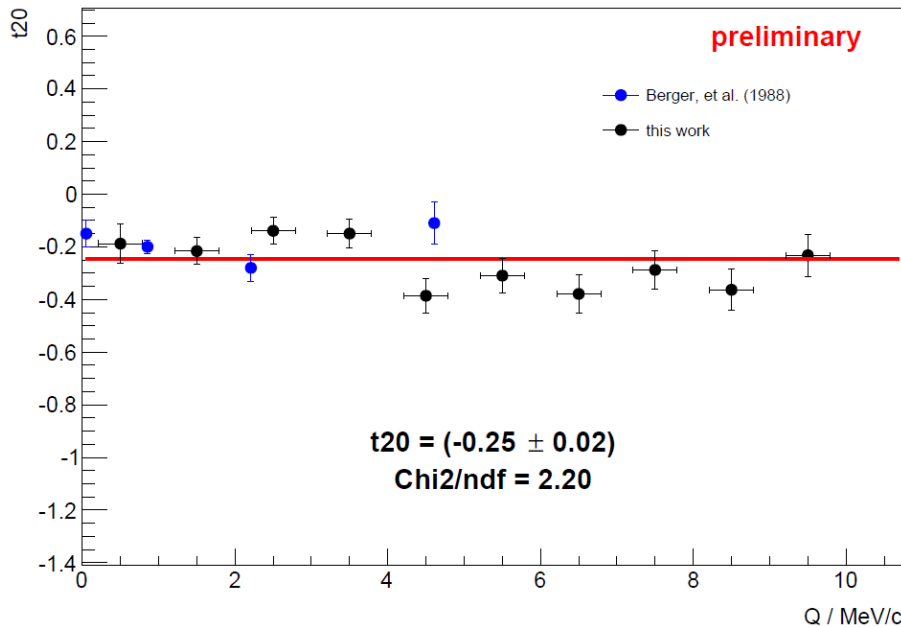


$$\frac{d\sigma_{\uparrow}(q, \varphi)}{dt} / \frac{d\sigma_0(q, \varphi)}{dt} =$$

$$1 + \sqrt{3} p_z t_{11}(\vartheta) \cos(\varphi) - \frac{1}{2\sqrt{2}} p_{zz} t_{20}(\vartheta)$$

$$- \frac{\sqrt{3}}{2} p_{zz} t_{22}(\vartheta) \cos(2\varphi)$$

# Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$



$$T_{20} = \frac{2 \cdot \sqrt{2}}{p_{zz}} \cdot \frac{d\sigma_0 / d\Omega(\mathcal{G}) - d\sigma_{\uparrow} / d\Omega(\mathcal{G})}{d\sigma_0 / d\Omega(\mathcal{G})}$$

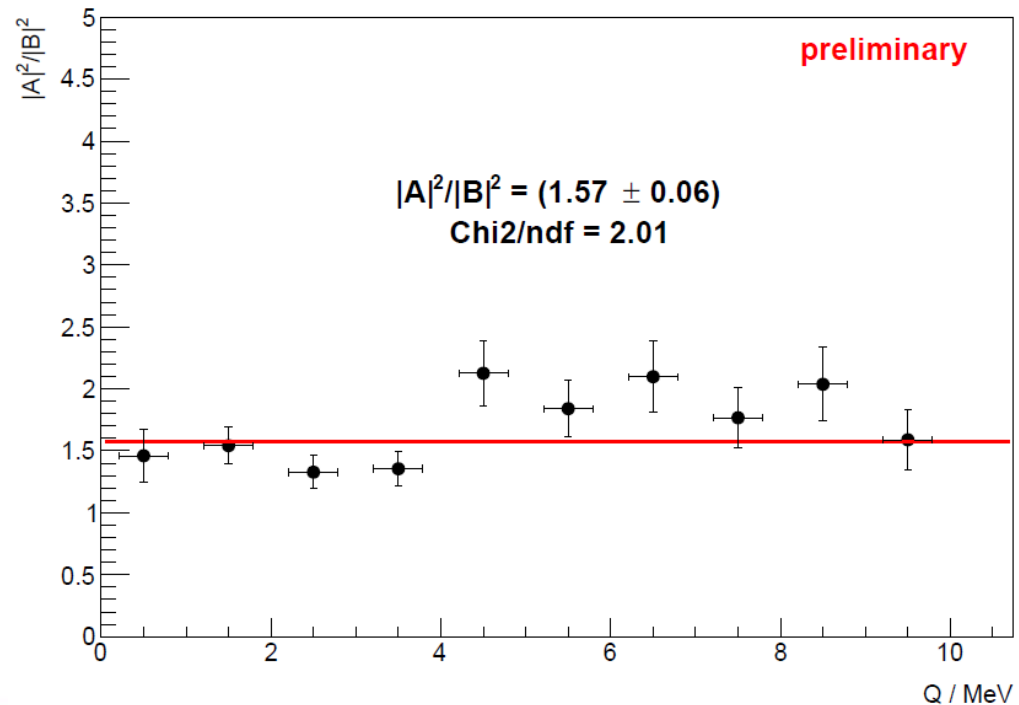
- Data indicate  $T_{20} = \text{const.}$  close to threshold
- $|T_{20}| \ll 1 \rightarrow |A|^2 / |B|^2 = O(1)$
- S-Wave amplitudes  $|A|^2$  and  $|B|^2$  are of similar size

# Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$

- Assumption:  $T_{20} = \text{const.} \rightarrow |A|^2/|B|^2 = \text{const.}$

$$T_{20} = \sqrt{2} \left[ \frac{|B|^2 - |A|^2}{|A|^2 + 2|B|^2} \right]$$

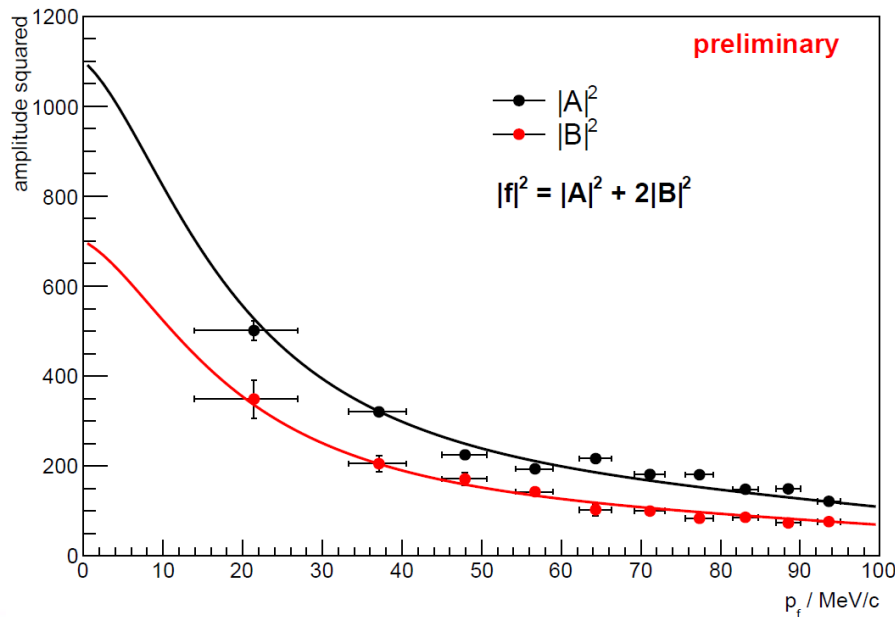
$$\rightarrow \frac{|A|^2}{|B|^2} = \frac{1 - \sqrt{2} \cdot T_{20}}{1 + T_{20} / \sqrt{2}}$$



# Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$

- Energy dependence of  $|f|^2$  known from „old“ unpolarized measurements

→  $|A|^2(p_f)$  and  $|B|^2(p_f)$  can be calculated



$$\frac{d\sigma}{d\Omega} = \frac{1}{3} \frac{p_\eta}{p_p} \left[ |A|^2 + 2|B|^2 \right]$$

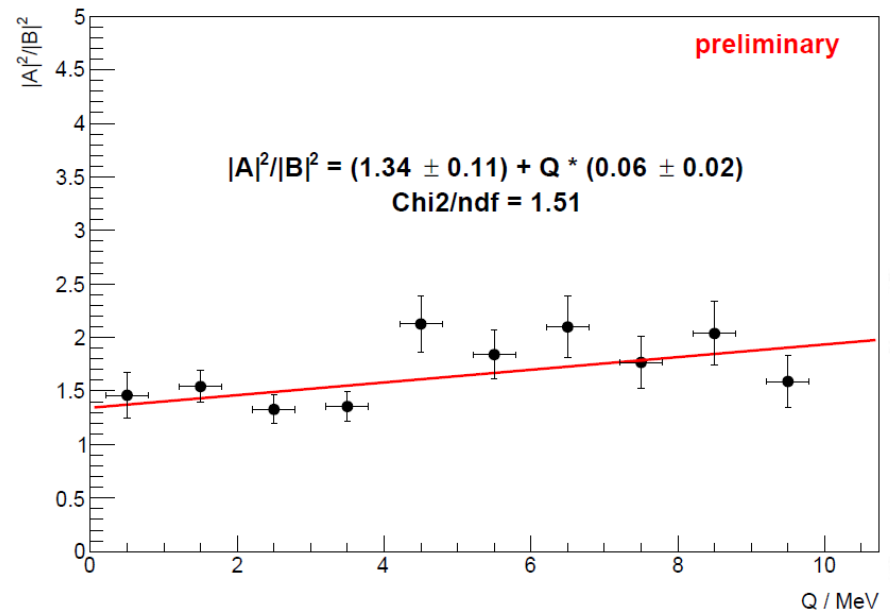
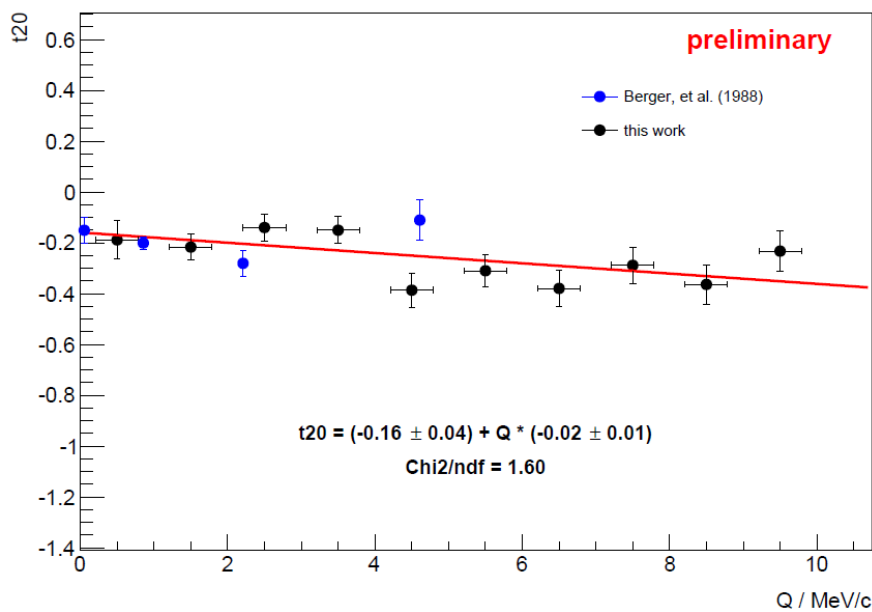
$$|A|^2 = \frac{p_p}{p_\eta} (1 - \sqrt{2}T_{20}) \frac{d\sigma}{d\Omega}$$

$$|B|^2 = \frac{p_p}{p_\eta} \left( 1 + \frac{1}{\sqrt{2}} T_{20} \right) \frac{d\sigma}{d\Omega}$$



# Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$

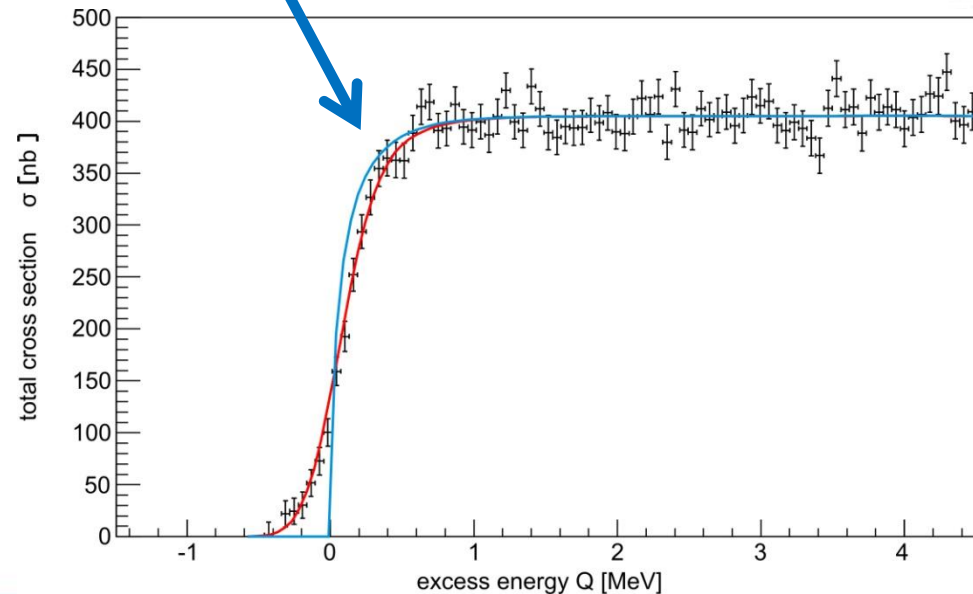
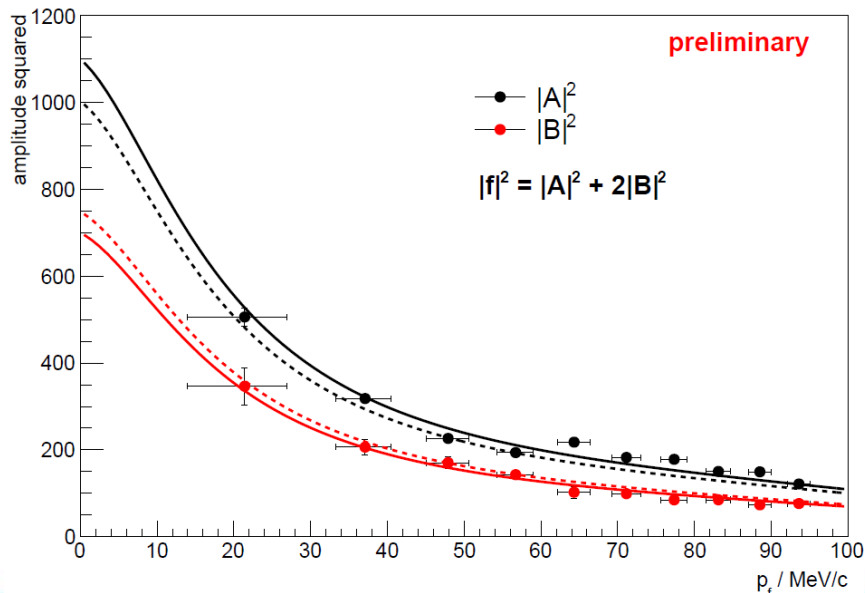
- Allow for an energy dependence of  $|A|^2/|B|^2$ :  
→ Test: Different energy dependence of  $|A|^2(p_f)$  and  $|B|^2(p_f)$  ?



$$\frac{d\sigma}{d\Omega} = \frac{1}{3} \frac{p_\eta}{p_p} \left[ |A|^2 + 2|B|^2 \right] \quad \left| \frac{|A|^2}{|B|^2} = m \cdot Q + n \right.$$

# Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$

- No significant different energy dependence of  $|A|^2$  and  $|B|^2$
- Remarkable excitation function of  $d+p \rightarrow {}^3\text{He}+\eta$  still an indication for very strong FSI effect

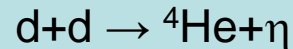
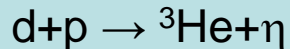
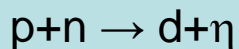


## Next Steps:

- Finalize data analysis
- Quantification of  $T_{20}$  and  $|A|^2/|B|^2$
- Estimation (or upper limits) for non-FSI effect
- Evaluation of effect on pole position or scattering length

In parallel:

- Analysis of new data on  $p+n \rightarrow d+\eta$  via  $p+d \rightarrow d+\eta+p_{\text{spec}}$
- Comparison of results from:



## Summary

- The  $\eta$ - $^3\text{He}$  system
  - exposes an unexpected strong final state interaction
  - is a good candidate for a bound meson-nucleus state (strong interaction)
- Preliminary tensor polarized data support the strong FSI interpretation
  - Possible spin-dependent effect gives only a minor contribution to the energy dependence of the production amplitude close to threshold
- New data on the  $d\eta$  system will allow for further investigations on the pole positions as function of the nucleus mass

## What else?

- There is need for further theoretical studies
  - on the extraction of FSI parameters from data
  - on the description of the production process: Two-Step Model etc.
- COSY offers nice possibilities for studies on bound states, but will stop soon the hadron physics program
  - What should be measured now?
  - What can we learn from possible new data?